Collapse and expansion of a Bose-Fermi mixture

Michele Modugno

INFM - LENS - Dipartimento di Fisica Università di Firenze, Italy

INFM Research and Development Center on Bose-Einstein condensation Bose-Einstein condensation @ Trento, Italy



Motivations

• Experiments with ⁴⁰K-⁸⁷Rb mixtures @ LENS:



Attractive Fermi-Bose interaction

[A. Simoni, F. Ferlaino, G. Roati, G. Modugno, and M. Inguscio, cond-mat/0301159]

 $\mathbf{a_{bf}} \simeq -410 \pm 80 \ \mathbf{a_0} \qquad (9/2, 9/2) \times (2, 2)$

Effects of a_{bf} on the expansion of bosons [G. Roati, F. Riboli, G. Modugno, and M. Inguscio, Phys. Rev. Lett. 89, 150403 (2002)]



Boson aspect ratio during the expansion

rightarrow Attractive interaction rightarrow Collapse

[G. Modugno, G. Roati, F. Riboli, G. Modugno, F. Ferlaino, R. J. Brecha, and M. Inguscio, Science **297**, 2240 (2002)]



Evolution of the number of atoms in the mixture during the evaporative cooling.

Outline of the talk

Mean-field analysis of:

X stability and collapse of the mixture

with: Exp. group @ LENS

```
rac{a}{b}f estimate of a_{bf}
```

X effects of the interaction on the expansion of fermions and bosons with: H. Hu (Post-Doc @ ICTP, Trieste) and X.-J. Liu (Post-Doc @ LENS) cond-mat/0301182

scaling approach

System geometry

Trapping potentials:

$$V_b(\mathbf{x}) = \frac{1}{2} m_b \omega_{b\perp}^2 \left[(x^2 + y^2) + \lambda^2 z^2 \right]$$

$$V_f(\mathbf{x}) = \frac{1}{2} m_f \omega_{f\perp}^2 \left[(x^2 + (y - y_0)^2) + \lambda^2 (z - z_0)^2 \right]$$

$$ightarrow \omega_{f\perp} = 2\pi imes 317$$
 Hz, $\lambda^{-1} \simeq 13.2$
 $\omega_{b\perp}$ a factor $\sqrt{m_b/m_f} \simeq 1.47$ smaller

× horizontal and vertical gravitational sag: $y_0 \simeq 3 \ \mu m$ and $z_0 \simeq 10 \ \mu m$



Typical geometry in the experiments @ LENS

A mean-field approach

GPE for bosons, Thomas-Fermi for fermions

[R. Roth, Phys. Rev. A 66, 013614 (2002): spherical symmetry, effective confinement]

$$\begin{cases} \left[-\frac{\hbar^2}{2m_b} \nabla^2 + V_b + g_{bb} n_b + g_{bf} n_f \right] \phi = \mu_b \phi \\ n_f = \frac{\left(2m_f\right)^{3/2}}{6\pi^2 \hbar^3} \left(\epsilon_f - V_f - g_{bf} n_b\right)^{3/2} \end{cases}$$



Ground-state of the mixture: density profiles along the vertical y direction.

Stability and collapse

– Typeset by $\mbox{Foil}{\rm T}_{\!E}\!{\rm X}$ –



 \sim Mean-field estimate for a_{bf}

$$a_{bf}^{mf} \simeq -395 \pm 15 \ a_0$$

in agreement with $a_{bf}^{coll} \simeq -410 \pm 80 \ a_0$

Beyond mean-field.....

[A. Albus, F. Illuminati, M. Wilkens, cond-mat/0211060]: exchange-correlations $\sim +a_{bf}^2$

Expansion

BEC: Thomas-Fermi hydrodynamic equations

$$\begin{cases} \frac{\partial n_b}{\partial t} + \nabla \cdot (n_b \mathbf{v}_b) = 0\\ m_b \frac{\partial \mathbf{v}_b}{\partial t} + \nabla \left(\frac{1}{2}m_b \mathbf{v}_b^2 + V_b + g_{bb}n_b + g_{bf}n_f\right) = 0 \end{cases}$$

Fermi gas: Boltzmann-Vlasov kinetic equation

$$\frac{\partial f}{\partial t} + \mathbf{v}_f \cdot \frac{\partial f}{\partial \mathbf{r}} - \frac{1}{m_f} \frac{\partial V_f}{\partial \mathbf{r}} \cdot \frac{\partial f}{\partial \mathbf{v}_f} - \frac{\mathbf{g_{bf}}}{m_f} \frac{\partial n_b}{\partial \mathbf{r}} \cdot \frac{\partial f}{\partial \mathbf{v}_f} = 0$$

Ground state: Thomas-Fermi approximation

$$\begin{cases} V_{b}(\mathbf{r}) + g_{bb}n_{b}^{0}(\mathbf{r}) + g_{bf}n_{f}^{0}(\mathbf{r}) = \mu_{b} \\ \frac{\hbar^{2}}{2m_{f}} \left(6\pi^{2}n_{f}^{0}(\mathbf{r})\right)^{2/3} + V_{f}(\mathbf{r}) + g_{bf}n_{b}^{0}(\mathbf{r}) = \mu_{f} \end{cases}$$

Simplifying assumption: *concentric* configuration

$$V_{b,f}\left(\mathbf{r}\right) = \frac{1}{2}m_{b,f}\omega_{\perp b,f}^{2}\left(\rho^{2} + \lambda^{2}z^{2}\right)$$

 $g_{bf} = 0$: simple scaling solution

for the condensate:

$$\begin{cases} n_b \left(\mathbf{r}, t\right) = \frac{1}{\prod_j b_j \left(t\right)} n_b^0 \left(\frac{r_i}{b_i(t)}\right) \\ v_{bi} \left(\mathbf{r}, t\right) = \frac{1}{\frac{b_i(t)}{b_i(t)}} \frac{db_i(t)}{dt} r_i \end{cases}$$

X scaling parameters:

$$\ddot{b}_{i}(t) + \omega_{bi}^{2}(t)b_{i}(t) - \frac{\omega_{bi}^{2}(0)}{b_{i}(t)\prod_{j}b_{j}(t)} = 0$$

and for the degenerate Fermi gas:

$$\begin{cases} f\left(\mathbf{r}, \mathbf{v}_{f}, t\right) = f_{0}\left(\frac{r_{i}}{\gamma_{i}(t)}\mathbf{V}(\mathbf{r}, t)\right)\\\\V_{i}(\mathbf{r}, t) = \gamma_{i}(t)v_{fi} - \frac{d\gamma_{i}(t)}{dt}r_{i} \end{cases}$$

X scaling parameters:

$$\ddot{\gamma}_i(t) + \omega_{fi}^2(t)\gamma_i(t) - \frac{\omega_{fi}^2(0)}{\gamma_i^3(t)} = 0$$

General case $oldsymbol{g}_{bf}
eq 0$

We require the scaling ansatz to be valid on average

[D. Guéry-Odelin, Phys. Rev. A 66, 033613 (2002)]: Collective oscillations in a classical gas

[C. Menotti, P. Pedri and S. Stringari, Phys. Rev. Lett. **89**, 250402 (2002)]: Expansion of an interacting Fermi gas

[X.-J. Liu and H. Hu Phys. Rev. A **67**, 023613 (2003)]: Collective oscillations in a Bose-Fermi mixture

The proper shapes of the density distributions do not enter directly the equations

only the knowledge of the initial equilibrium density distribution is required!

– Typeset by $\mbox{FoilT}_{\!E\!} X$ –

Bosons:

$$\begin{split} \ddot{b}_{i}(t) + \omega_{bi}^{2}(t)b_{i}(t) - \frac{\omega_{bi}^{2}(0)}{b_{i}(t)\prod_{j}b_{j}(t)} - \frac{g_{bf}}{m_{b}N_{b}\langle r_{i}^{2}\rangle_{b}} \frac{1}{b_{i}\prod_{j}b_{j}} \int d^{3}\mathbf{r} \frac{\partial n_{f}^{0}}{\partial r_{i}} r_{i}n_{b}^{0} \\ + \frac{g_{bf}}{m_{b}N_{b}\langle r_{i}^{2}\rangle_{b}} \frac{1}{b_{i}\prod_{j}b_{j}} \int d^{3}\mathbf{r} \frac{\partial n_{f}^{0}}{\partial r_{i}} r_{i}n_{b}^{0}(\frac{\gamma_{i}}{b_{i}}r_{i}) = 0 \end{split}$$

Fermions:

$$\begin{split} \ddot{\gamma}_{i}(t) + \omega_{fi}^{2}(t)\gamma_{i}(t) - \frac{\omega_{fi}^{2}(0)}{\gamma_{i}^{3}(t)} - \frac{g_{bf}}{m_{f}N_{f}\left\langle r_{i}^{2}\right\rangle_{f}} \frac{1}{\gamma_{i}^{3}} \int \! d^{3}\mathbf{r} \frac{\partial n_{b}^{0}}{\partial r_{i}} r_{i} n_{f}^{0} \\ + \frac{g_{bf}}{m_{f}N_{f}\left\langle r_{i}^{2}\right\rangle_{f}} \frac{1}{\gamma_{i}\prod_{j}\gamma_{j}} \int \! d^{3}\mathbf{r} \frac{\partial n_{b}^{0}}{\partial r_{i}} r_{i} n_{f}^{0}(\frac{b_{i}}{\gamma_{i}}r_{i}) = 0 \end{split}$$

– Typeset by $\mbox{FoilT}_{\!E\!X}$ –

static effect on the ground state profiles:

attractive interaction: both the densities are remarkably enhanced within the overlap region.

 \implies tighter confinement: if considered alone would lead to a faster expansion for both species.

dynamical effect during the early stages of the expansion: with the attractive interaction both species will feel a *running confinement* which reduces their expansion rate. The Due to statistics the Fermi distribution is wide even for small N: we expand n_f^0 around the center \longrightarrow for the bosons:

$$\ddot{b}_{i}(t) - \frac{\omega_{bi}^{2}(1+A)}{b_{i}(t)\prod_{j}b_{j}(t)} + \frac{A\omega_{bi}^{2}b_{i}(t)}{\gamma_{i}^{2}(t)\prod_{j}\gamma_{j}(t)} \approx 0$$
$$A \approx \frac{g_{bf}}{m_{b}N_{b}\langle r_{i}^{2}\rangle_{b}\omega_{bi}^{2}} \int d^{3}\mathbf{r} \frac{\partial n_{f}^{0}(\mathbf{r})}{\partial r_{i}} r_{i}n_{b}^{0}(\mathbf{r})$$

Static effect on bosons:

$$A \approx 40\% \longrightarrow \Delta \omega / \omega \approx 20\%$$

rightarrow two times larger than that used in the experimental fit (10%)



Boson aspect ratio during the expansion ($N_b = 2 \times 10^4$, $N_f = 10^4$).

we expect the "dynamical effect" to give important corrections: the Fermi-Bose interaction during the early stages of the expansion is **not negligible!**

$$b_{\perp}(\tau) = \sqrt{1 + (1 + \delta)^2 \tau^2}, \qquad b_z(\tau) \approx 1$$
$$\gamma_{\perp}(\tau) \approx \sqrt{1 + \beta^2 \tau^2}, \qquad \gamma_z(\tau) \approx 1$$

 $<\!\!\!>$ the *net* increase δ of the boson trapping frequencies is ($\beta \equiv \omega_f/\omega_b$)

$$\delta = \beta \left[\frac{\left(\beta^4 + 4A(1+A)\right)^{1/2} - \beta^2}{2A} \right]^{1/2} - 1$$

 $<\!\!\!> \delta = 12\%$, in good agreement with the exp. fitting value of 10%

Expansion of ⁸⁷Rb: the static effect (+20%) dominates over the dynamical one (-8%)

Full solution of the scaling equations:

X Expansion of ⁴⁰K: the dynamical effect always dominates and the aspect ratio is less than that of a pure Fermi gas.

X Dependence of the aspect ratio on the scattering lenght

X Dependence of the aspect ratio on the trap geometry

Summary

Mean-field analysis of stability and collapse of a Bose-Fermi system: 40K-87Rb mixtures @ LENS:

× estimate of $a_{bf}^{mf} \simeq -395 \pm 15 \ a_0$

Scaling approach to the expansion of the mixture:

X increase of the boson aspect ratioX reduction on the fermion aspect ratio

with respect to the noninteracting case.