

Hydrodynamics in a Degenerate Strongly-Interacting Fermi Gas



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Atom Cooling and Trapping

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Sponsored by: DOE, NSF, ARO, NASA

Outline



- Strongly-Interacting Fermi Gases
- All-Optical Cooling Method
- Hydrodynamic Expansion
- Trapped Atom Hydrodynamics – Fermi superfluid?
- Conclusions

Strongly-Interacting Fermi Gases

Mimic Exotic Systems in Nature



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A BUNCH OF DEGENERATES

A degenerate gas of fermions occurs in diverse situations, as described below:

■ Superconductors:

The electrons are degenerate and form loosely correlated Cooper pairs, which produce the superconductivity. Something similar must happen in high-temperature superconductors, but that process remains a mystery.

■ Neutron stars: The refusal of neutrons (which are fermions) to occupy identical quantum states generates a repulsion that prevents the star from collapsing under its own immense weight.

A similar repulsion stabilizes the laboratory-made degenerate fermi gases against collapse.

- High-Temperature Superconductors
- Neutron Stars
- Strongly-Interacting Matter
- Quark-Gluon Plasma – Elliptic Flow

■ Quark-gluon plasma:

As created at the Relativistic Heavy Ion Collider at Brookhaven National Laboratory, the exploding cloud of free quarks (which are fermions) and gluons has properties similar to a gas of fermionic atoms released from the confines of a trap.

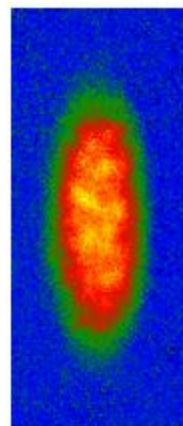


“Shape-shifting” Fermi Gas

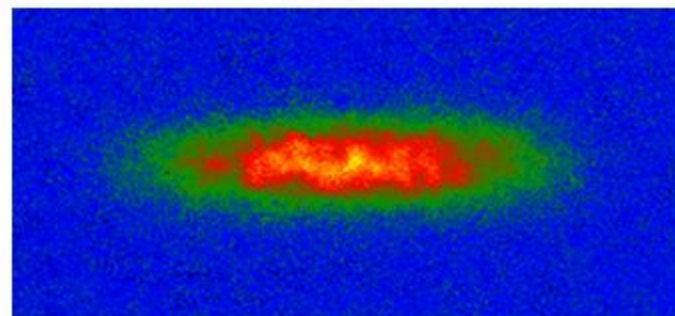
Trapped Gas:
0 ms



Released Gas:
after 2.1 ms



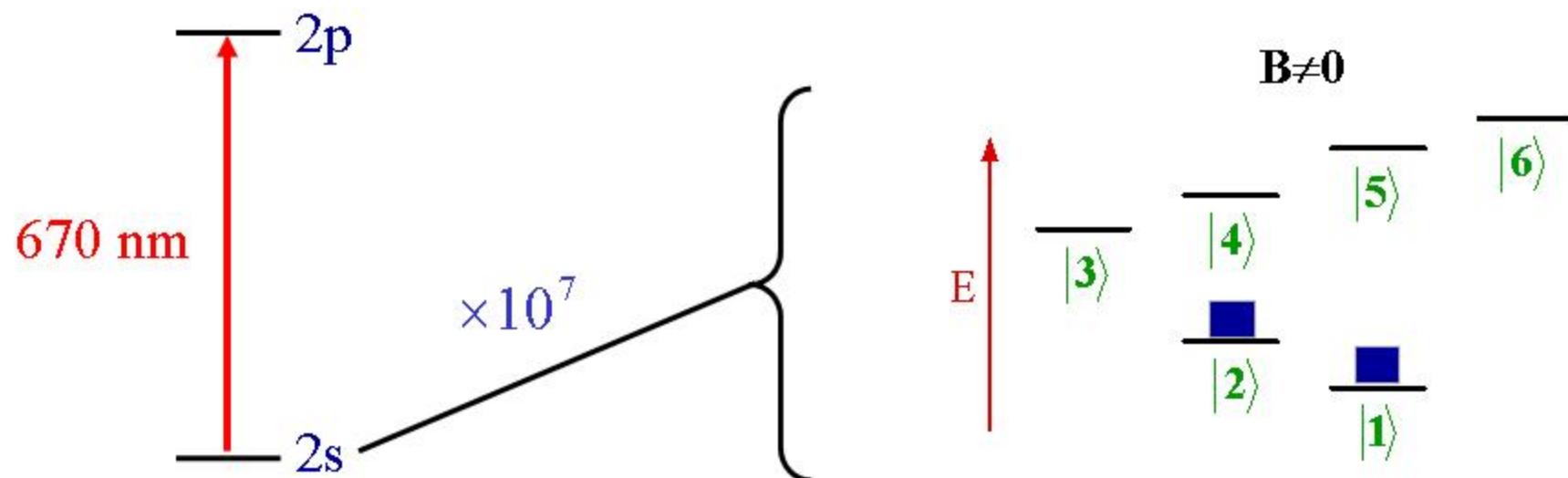
Released Gas: 12 ms



Fermi energy 700 nK
 $T < 35$ nK

Mixture of Spin Up/Down ${}^6\text{Li}$ Atoms

Ground State Hyperfine Structure in a Magnetic Field B



Ground State: $1s^2 2s^1$

Nuclear Spin: $I_N = 1$

Level 1 Spin $\frac{1}{2}$ Up

$$|1\rangle = |\uparrow\rangle = \left| -\frac{1}{2}, 1 \right\rangle$$

Level 2 Spin $\frac{1}{2}$ Down

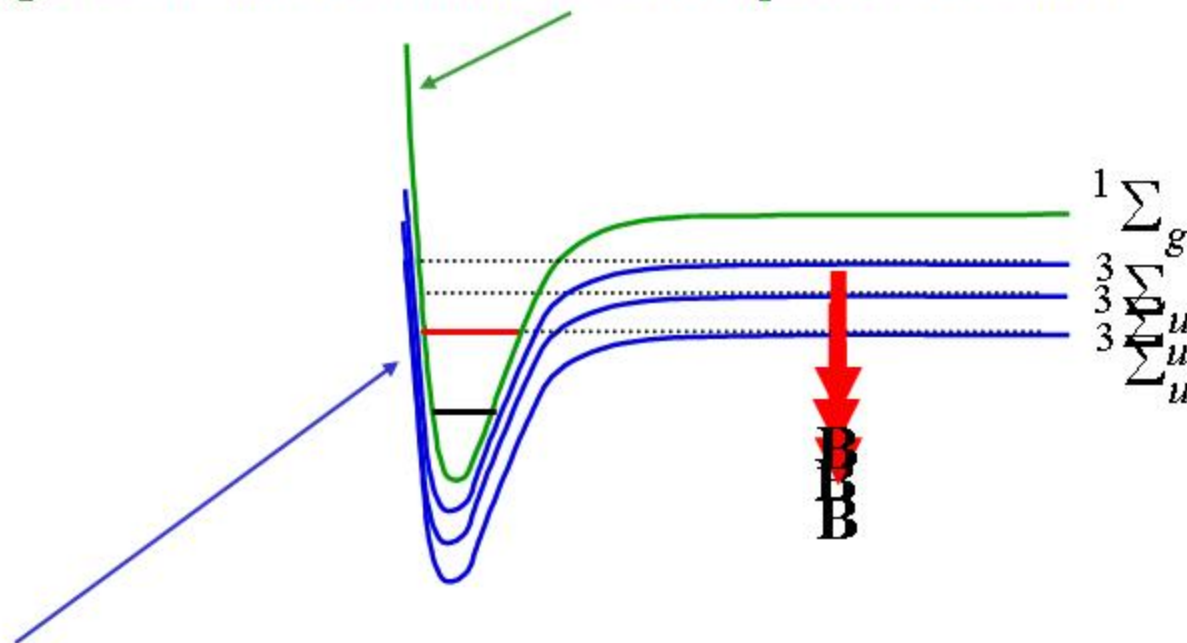
$$|2\rangle = |\downarrow\rangle = \left| -\frac{1}{2}, 0 \right\rangle$$

1,2 States High B-Field Seeking-Requires Optical Trap

Producing Strong Interactions with a Feshbach Resonance

Resonant Coupling between Colliding Atom Pair – Bound Molecular State

Singlet Diatomic Potential: Electron Spins Not Parallel



Triplet Diatomic Potential: Electron Spins Parallel

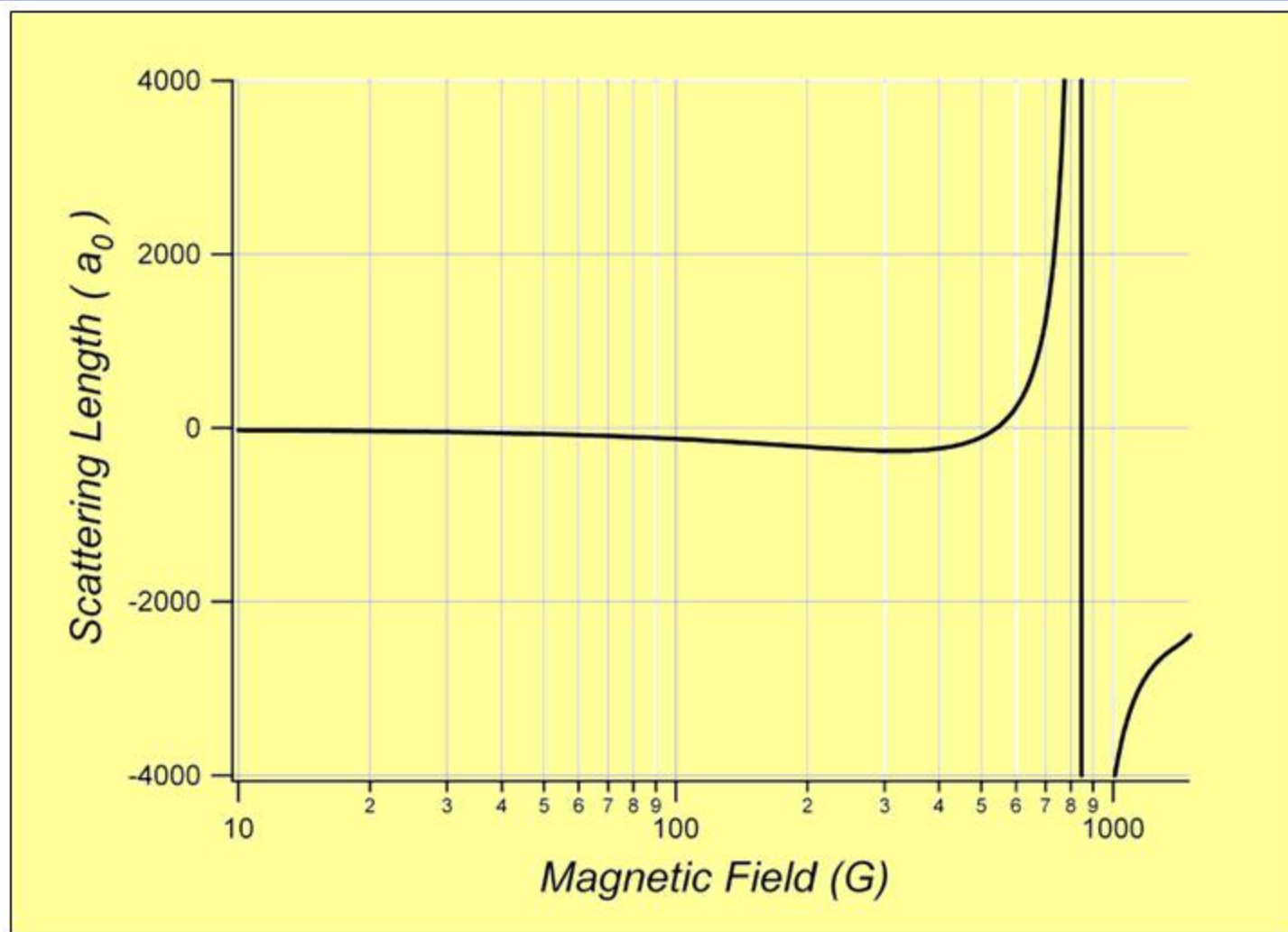
Zero Energy Scattering Length $a_s \rightarrow \pm\infty$

*Feshbach Resonance for 1-2 Mixture



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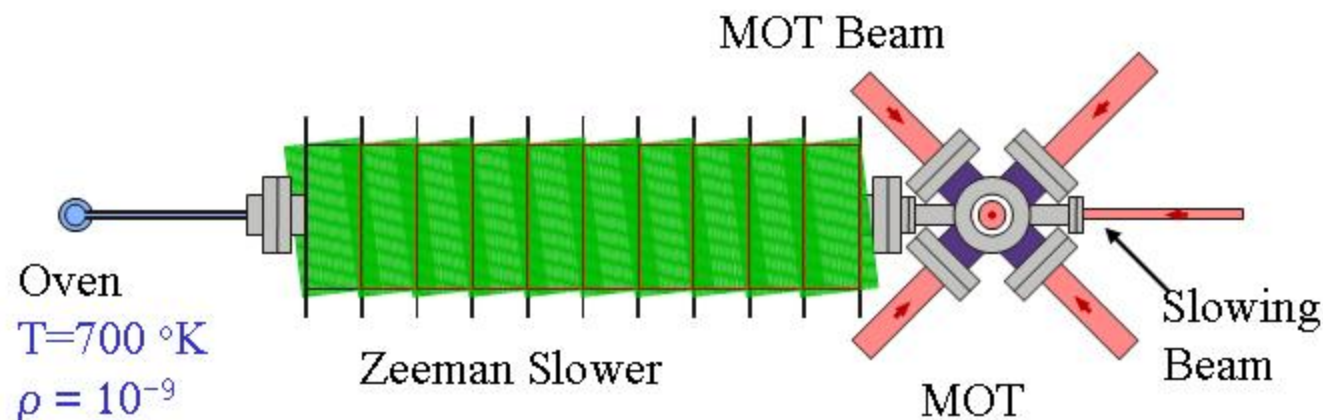
*M. Houbiers, *et al.*, Phys. Rev. A **57**, 1497 (1998)

Cooling Fermi Gases in Optical Traps



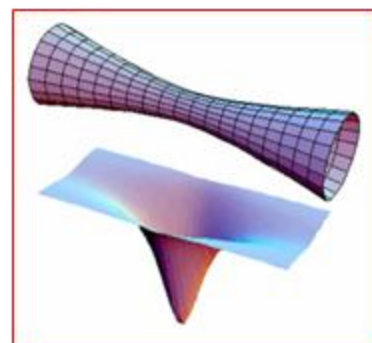
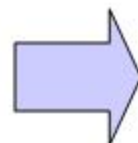
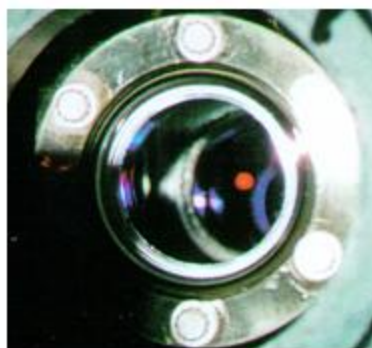
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$T < 100\text{ nK}$
 $\rho = 1$

$P = 10^{-11}\text{ Torr}$
 $T = 150\text{ }\mu\text{K}$
 $\rho = 10^{-5}$



Optical Trap

Optical Trap

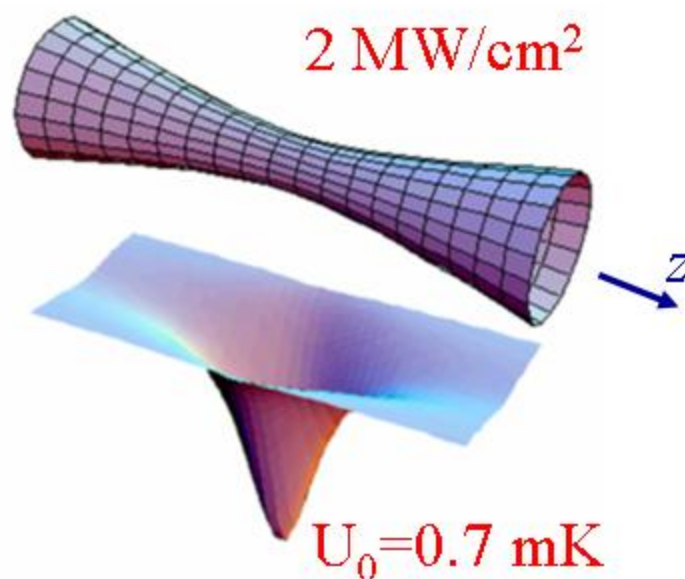


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Focused Gaussian Laser Beam

$$U = -\frac{1}{2} \alpha \overline{E_0^2} \frac{1}{1+(z/z_0)^2} e^{-2r^2/w_0^2}$$



Ultrastable CO₂ Laser Trap

- Stable Commercial Laser



- Typical Trap Parameters

$$P = 65 \text{ W} \quad \omega_0 = 47 \text{ } \mu\text{m}$$

$$z_0 = 0.7 \text{ mm}$$

$$I_0 = 2.0 \text{ MW/cm}^2 \quad U_0 = 0.7 \text{ mK}$$

$$\nu_r \cong 6.6 \text{ kHz} \quad \nu_z \cong 340 \text{ Hz}$$

- Negligible Optical Heating
 - Scattering Time $\cong 1/2$ hour
 - Optical Heating $\cong 18 \text{ pK/s}$

- Extremely Low Noise
 - Intensity Noise Heating

$$\Gamma^{-1} \geq 2.3 \times 10^4 \text{ sec}$$

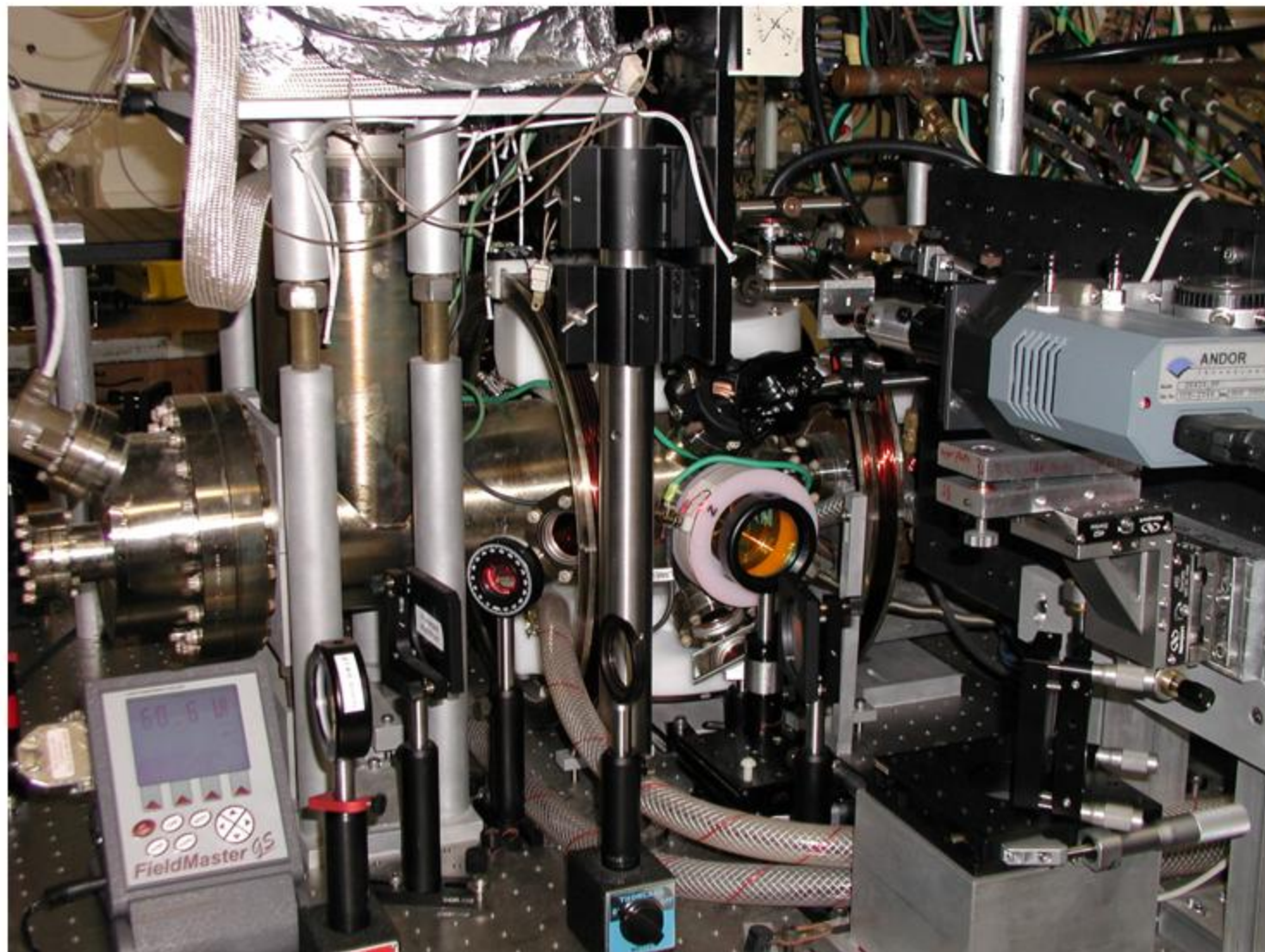
- Ultra-High Vacuum
 - Pressure: $< 10^{-11}$ Torr
 - Heating: $< 5 \text{ nK/sec}$
 - Lifetime: 400 sec

Experimental Apparatus



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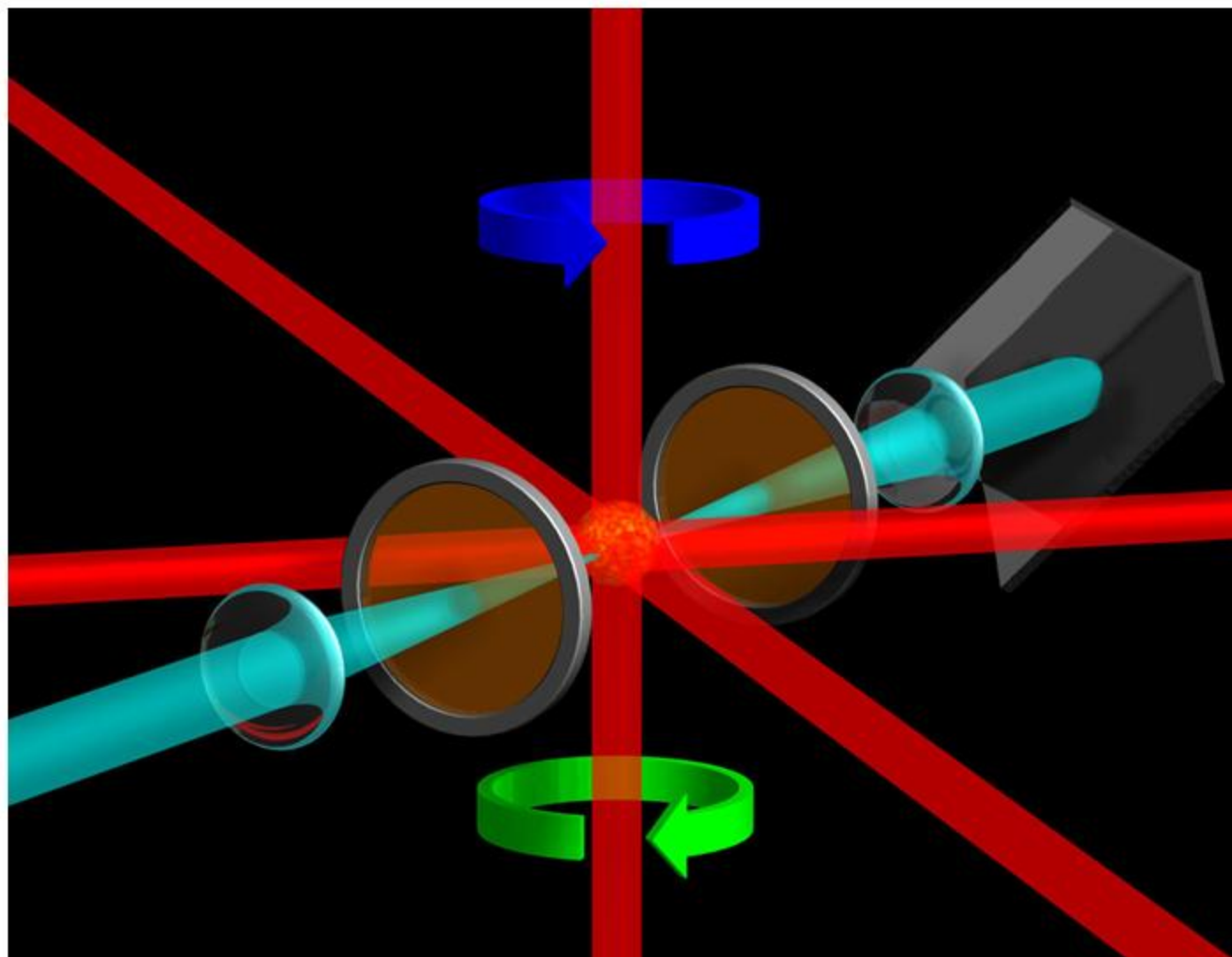


Optical Trap Loading



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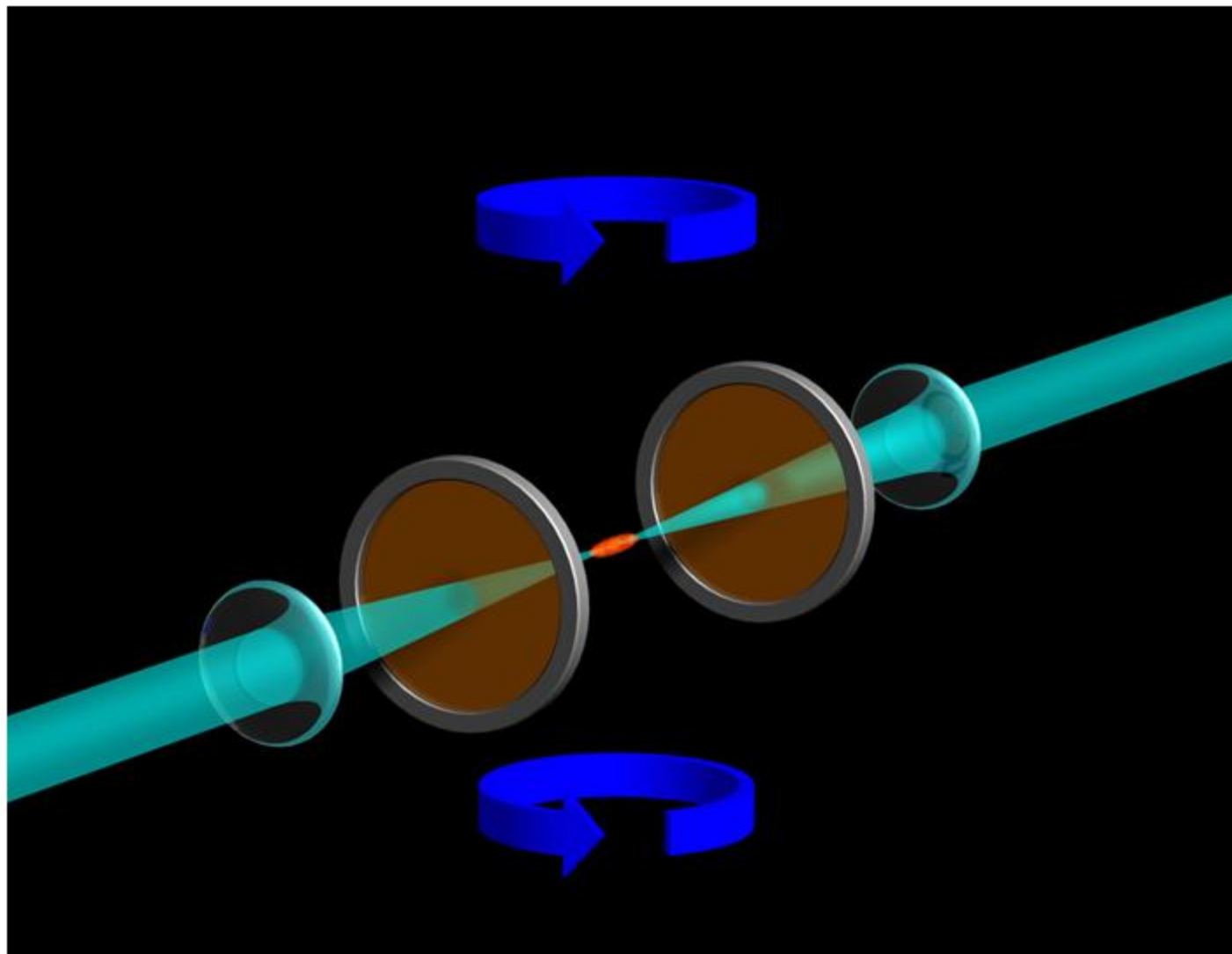


Forced Evaporation



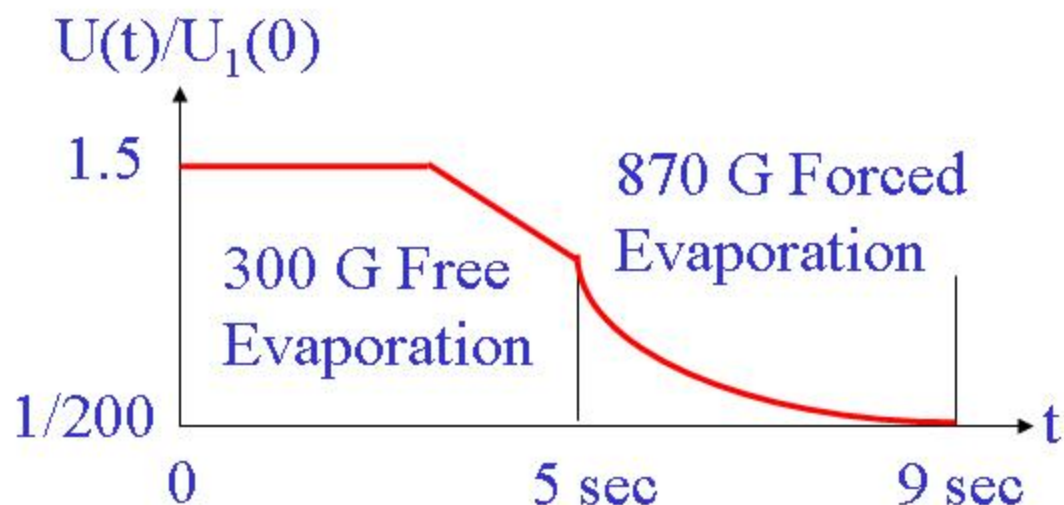
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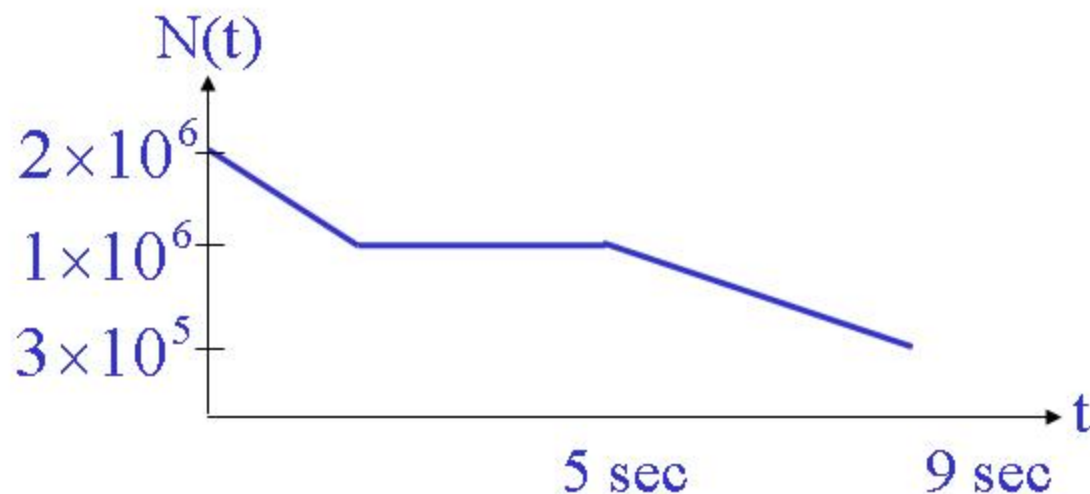


Timing Sequence for Evaporation



Scaling Laws:

$$\frac{\rho_f}{\rho_i} = \left(\frac{U_i}{U_f} \right)^{1.3} = 10^3$$



$$\frac{N_f}{N_i} = \left(\frac{U_f}{U_i} \right)^{\frac{3}{16}} = \frac{1}{2.7}$$

Cooling a Strongly-Interacting Fermi Gas

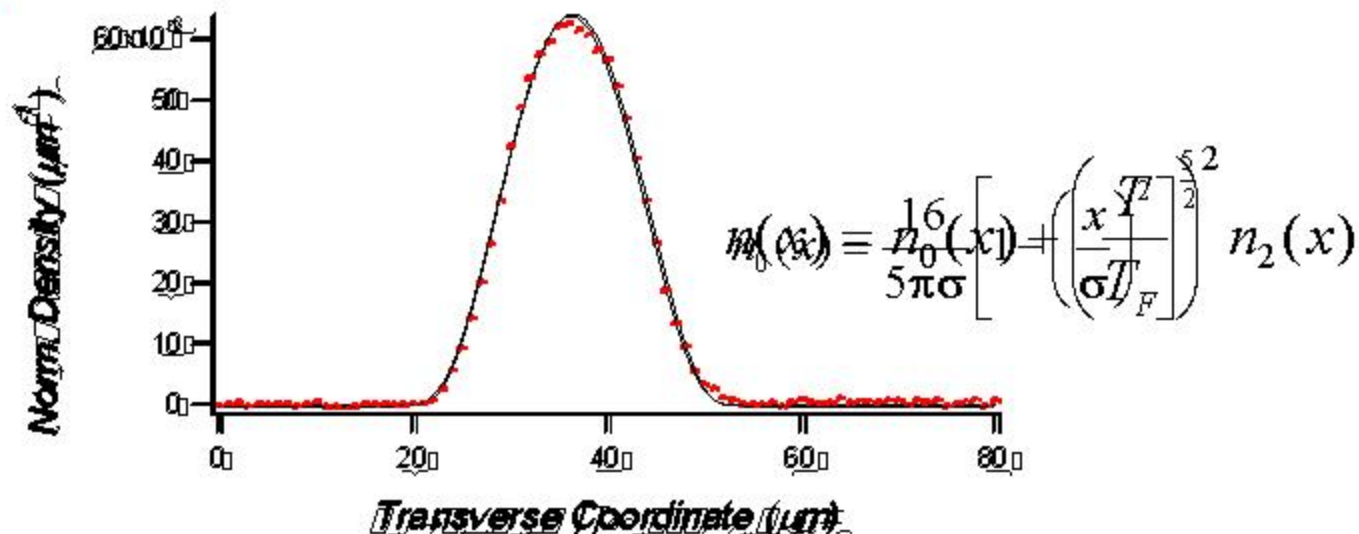
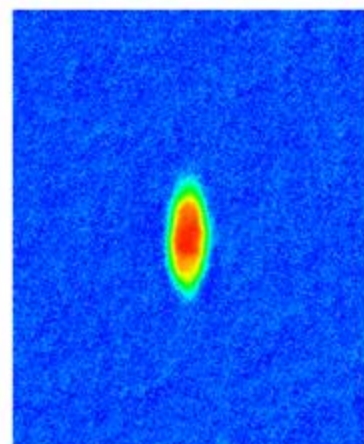


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Evaporate for 3.5 s at 910 G:

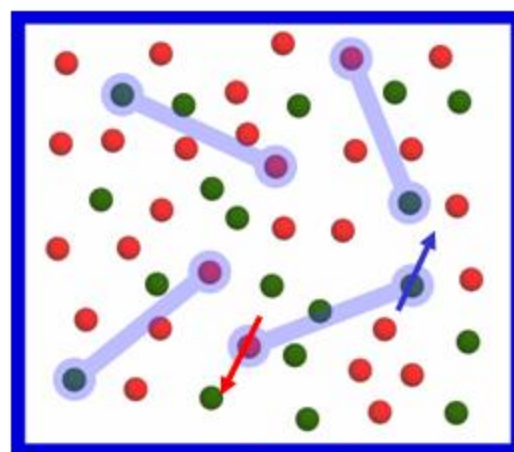
$t = 0.2$ ms after release



$T/T_F = 0.09$
 $T = 0.7 \mu\text{K}$ at full trap depth U_0
 $T = 50 \text{ nK}$ at $U_0/200$

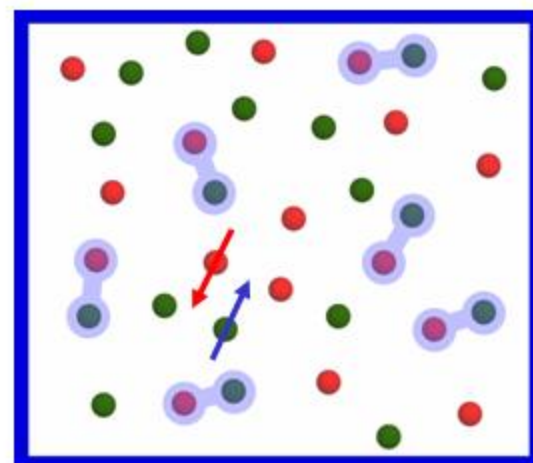
Superfluidity in Atomic Fermi Gases

- Magnetically tunable interactions via Feshbach Resonance
- Theory BCS Pairing ^6Li : Houbiers, et al., PRA **56**, 4864 (1997).



$$\eta_C \approx \exp\left(\frac{-L}{|a_s|}\right) \ll 1$$

$$T_C = \eta_C T_F$$



$$\eta_C \approx 0.2-0.5$$

- On Resonance: Super-High T_C Superconductivity

M. Holland, *et al.* Phys. Rev. Lett. **87**, 120406 (2001)

E. Timmermans, *et al.* Phys. Lett. A **285**, 228 (2001)

Y. Ohashi *et al.* cond-mat/0201262 (2002)

Prediction of Anisotropic Expansion*

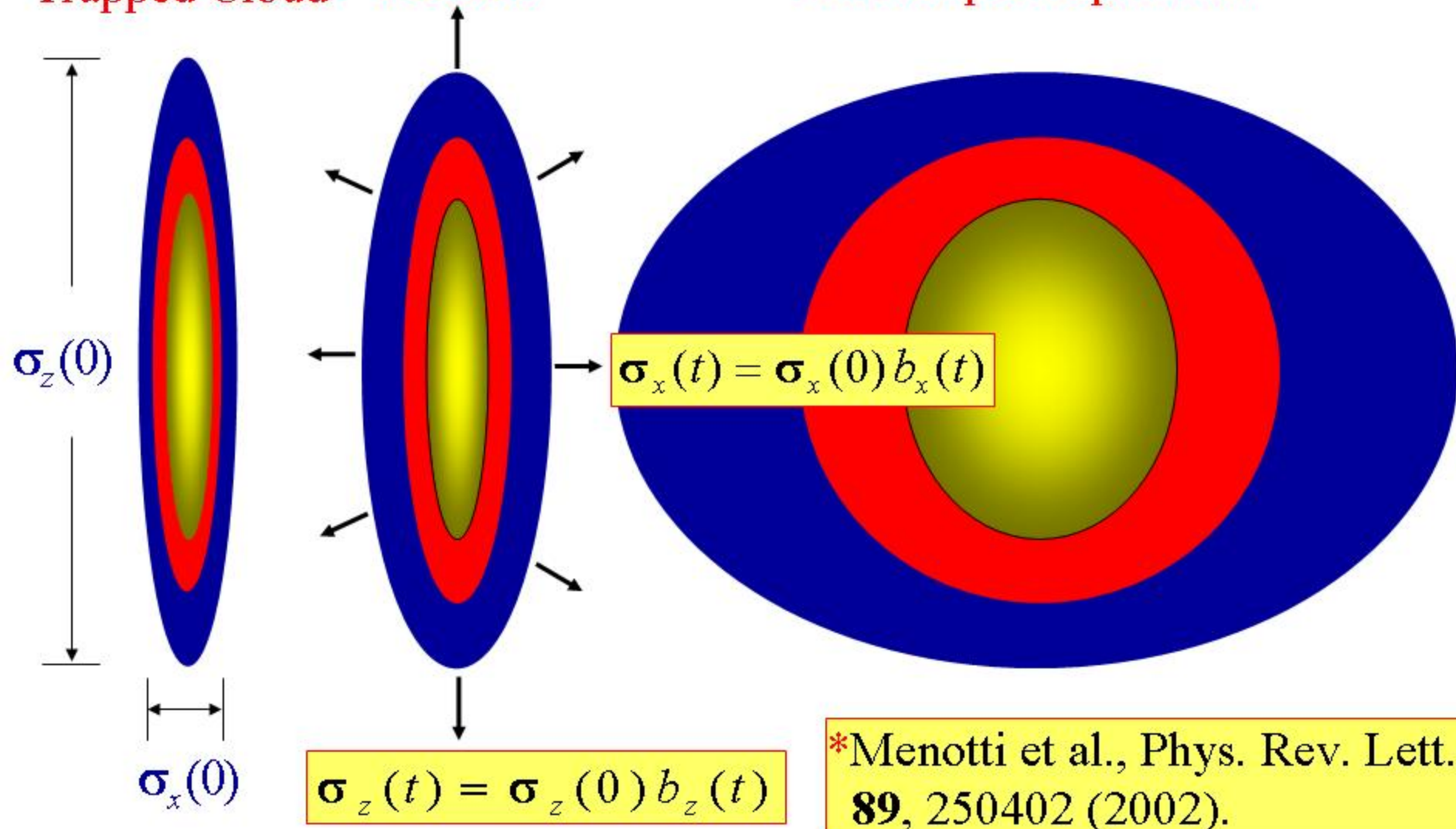


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Trapped Cloud Release

Anisotropic Expansion

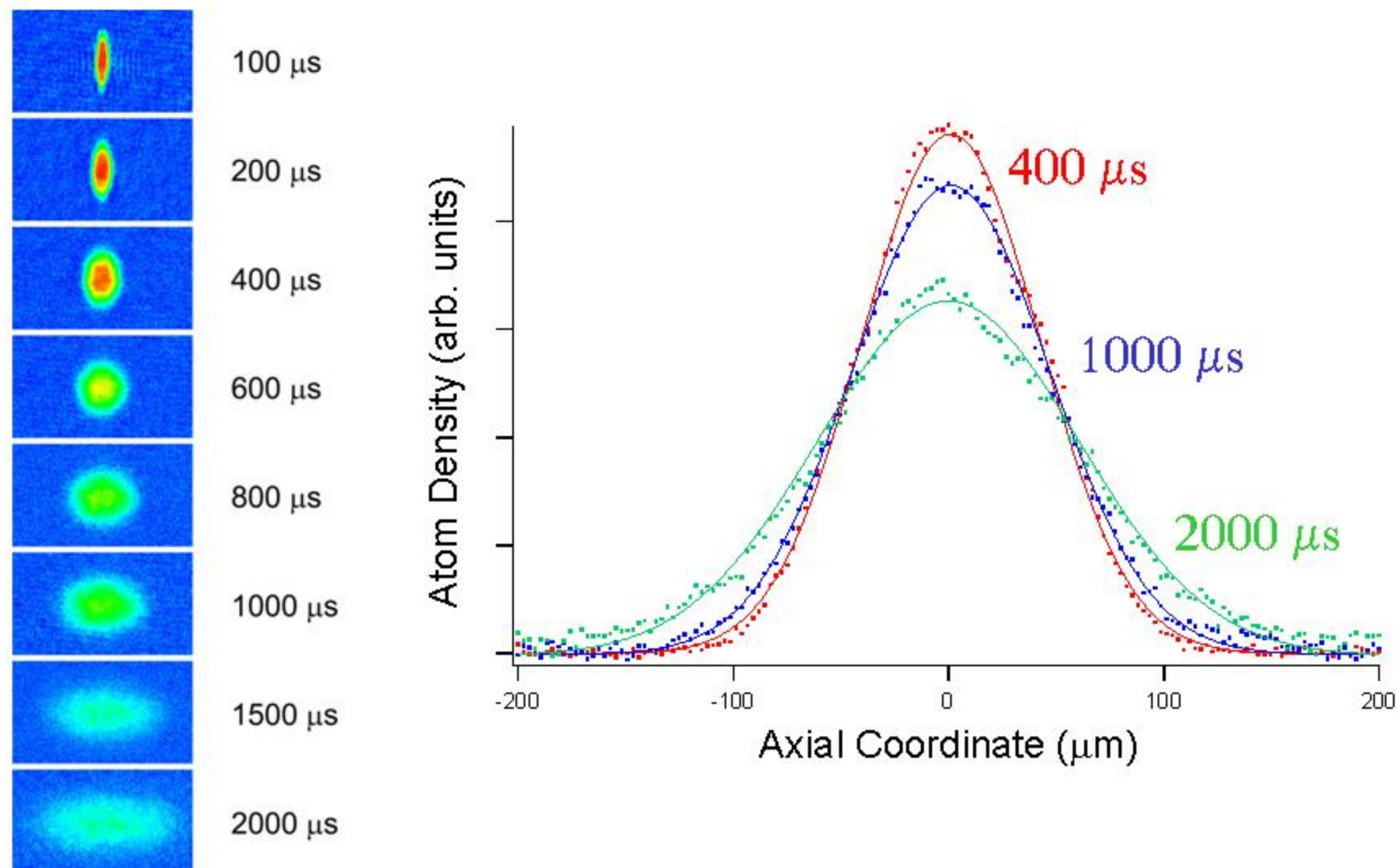


Axial Expansion at 910 G



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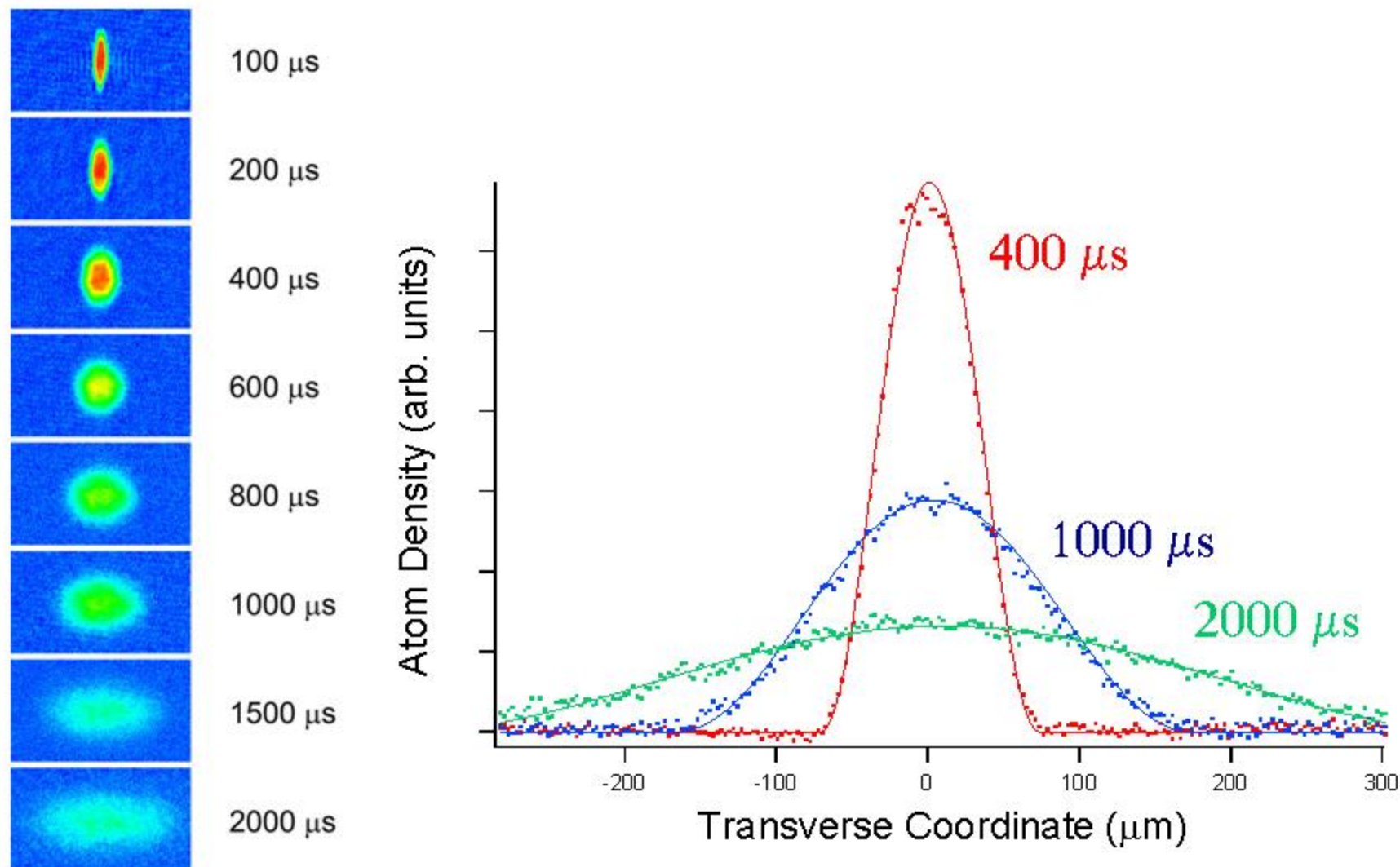


Transverse Expansion at 910 G



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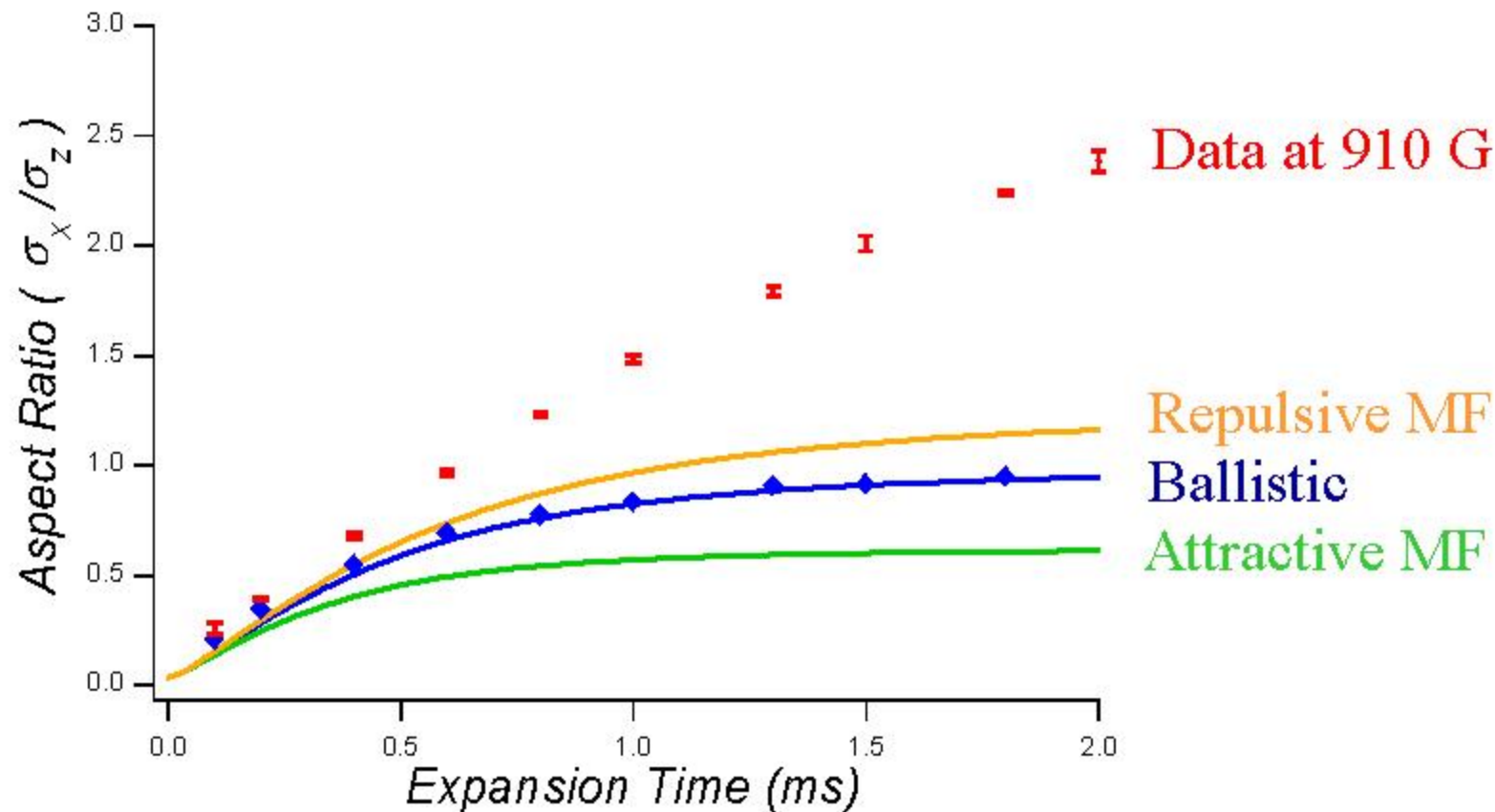


Hydrodynamic Expansion Data



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Superfluid Hydrodynamic Expansion



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Continuity Equation

$$\frac{\partial n}{\partial t} + \nabla \cdot (n \mathbf{u}) = 0$$

Irrotational

$$\nabla \times \mathbf{u} = 0$$

Euler's equation for a Superfluid

$$M \left(\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) = -\nabla \mu(n)$$

Chemical potential with mean field

$$\mu = (1 + \beta_S) \varepsilon_F(x)$$

$$M \frac{\partial \mathbf{u}}{\partial t} = -\nabla \left((1 + \beta_S) \varepsilon_F(x) + \frac{1}{2} M \mathbf{u}^2 \right)$$

Expansion of the Interacting Gas

The form of the density distribution does not change:

$$n(x, t) = \frac{1}{b_x(t)} n_0 \left(\frac{x}{b_x(t)} \right)$$

Good News:

The time-dependent scale factors obey simple differential equations:

$$\ddot{b}_i = \frac{\omega_i^2}{(b_x b_y b_z)^{2/3} b_i} \quad b_i(0) = 1, \quad \dot{b}_i(0) = 0$$

Bad News: Independent of β !

*Kagan et al., Phys. Rev. A **55**, R18 (1997).

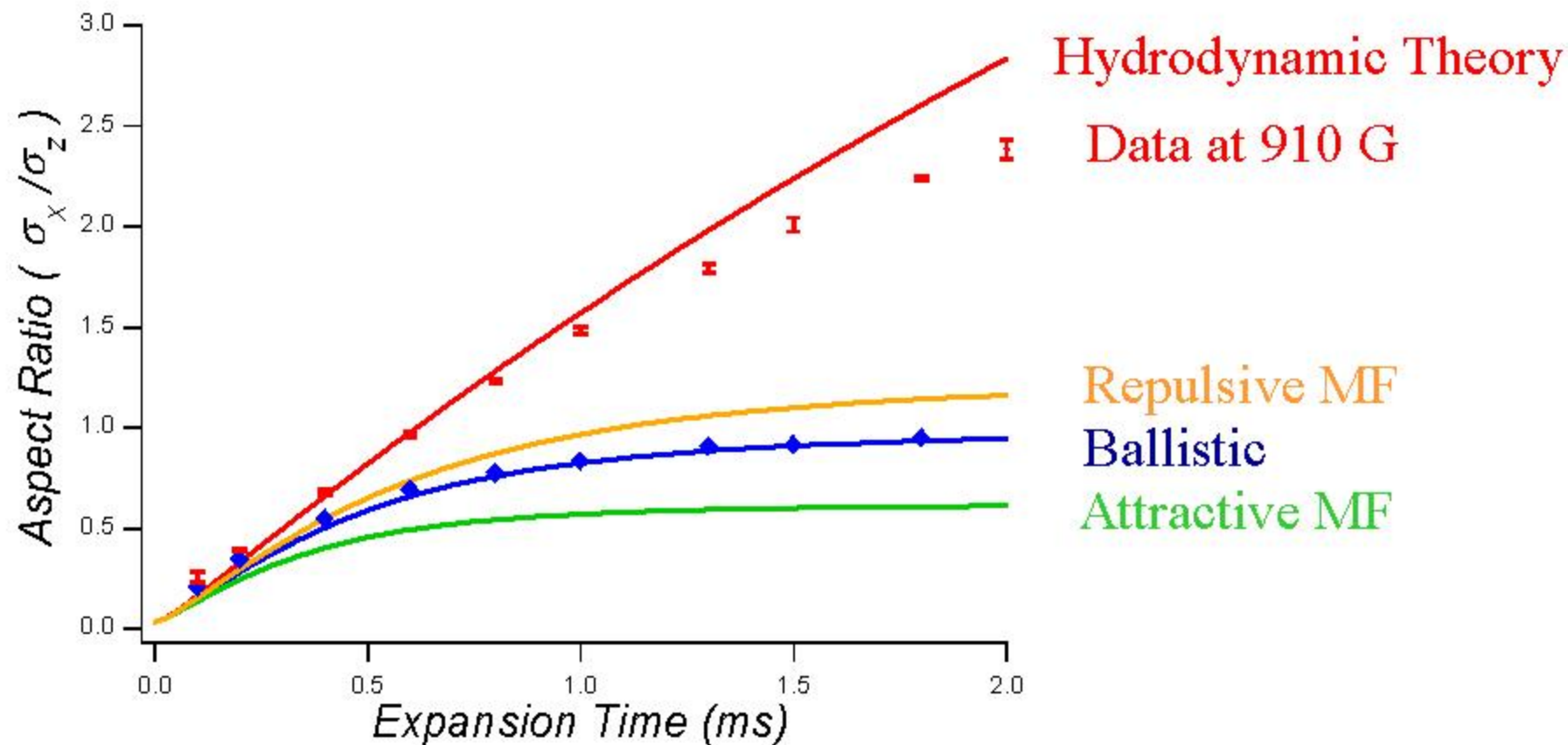
*Menotti et al., Phys. Rev. Lett. **89**, 250402 (2002).

Expansion Dynamics



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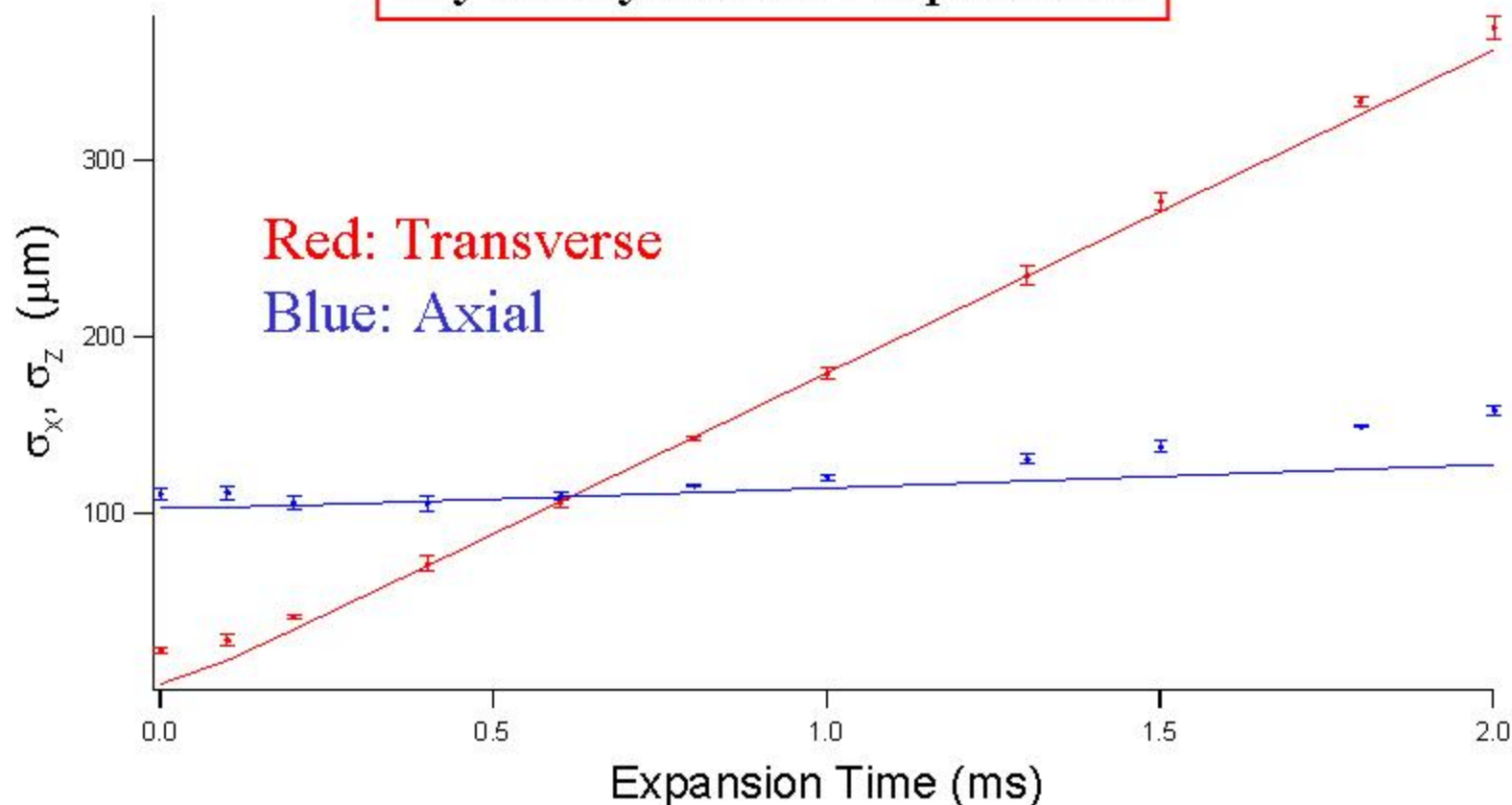
Transverse and Axial Widths vs Time



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Hydrodynamic Expansion



Collisional Hydrodynamic Expansion?



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Continuity Equation

$$\frac{\partial n}{\partial t} + \nabla \cdot (n \mathbf{u}) = 0$$

Irrotational

$$\nabla \times \mathbf{u} = 0$$

Euler's equation for a normal fluid

$$M \left(\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} \right) = - \frac{\nabla P}{n}$$

Pressure

$$P(x) = \frac{2}{3} \frac{\langle E \rangle}{V} = \frac{2}{5} (1 + \beta_N) \varepsilon_F(x) n(x)$$

$$M \frac{\partial \mathbf{u}}{\partial t} = - \nabla \left((1 + \beta_N) \varepsilon_F(x) + \frac{1}{2} M \mathbf{u}^2 \right)$$

Same equation!



Is collisional dynamics possible ?

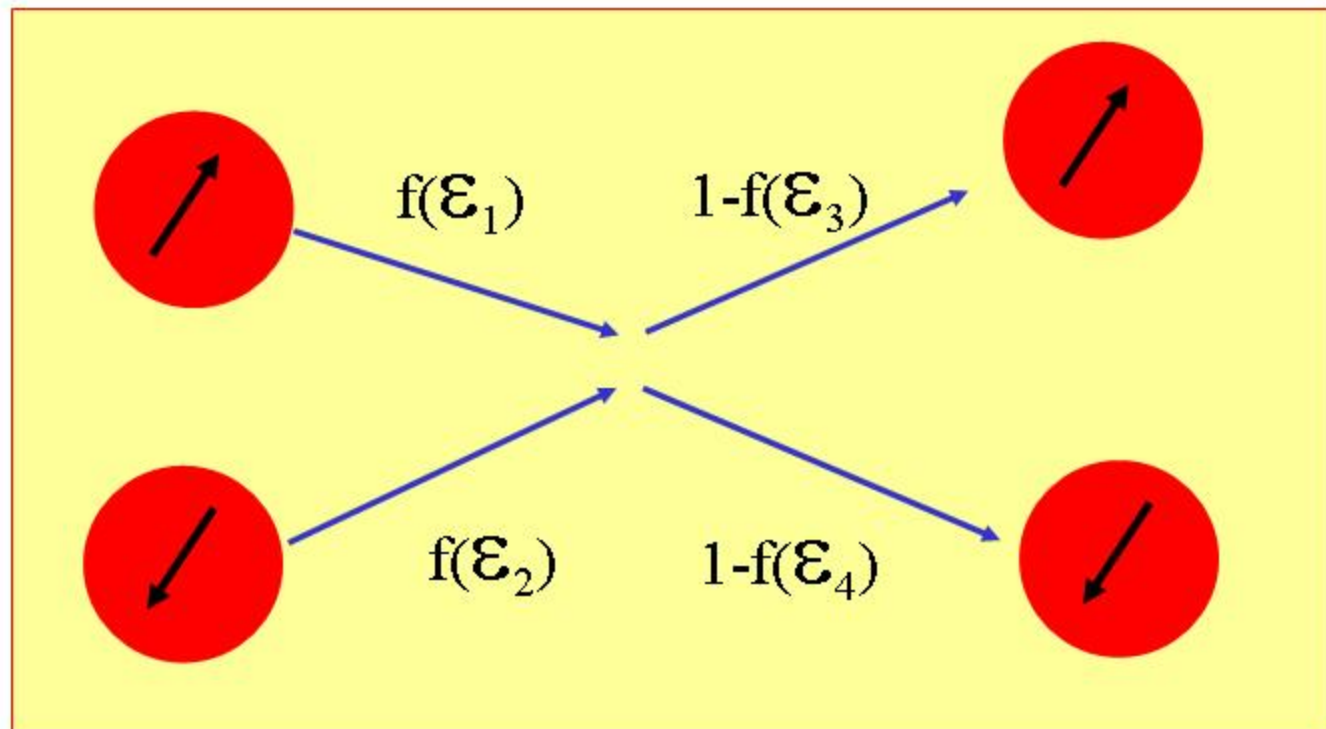
Can Pauli Blocking Suppress Collisions?



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Collision **cannot** occur if the final state is occupied:



$$\Gamma_{\text{coll}} = \gamma \left(\frac{T}{T_F} \right)^2$$

→ 0

as $T \rightarrow 0$

$f(\epsilon)$ is the occupation probability

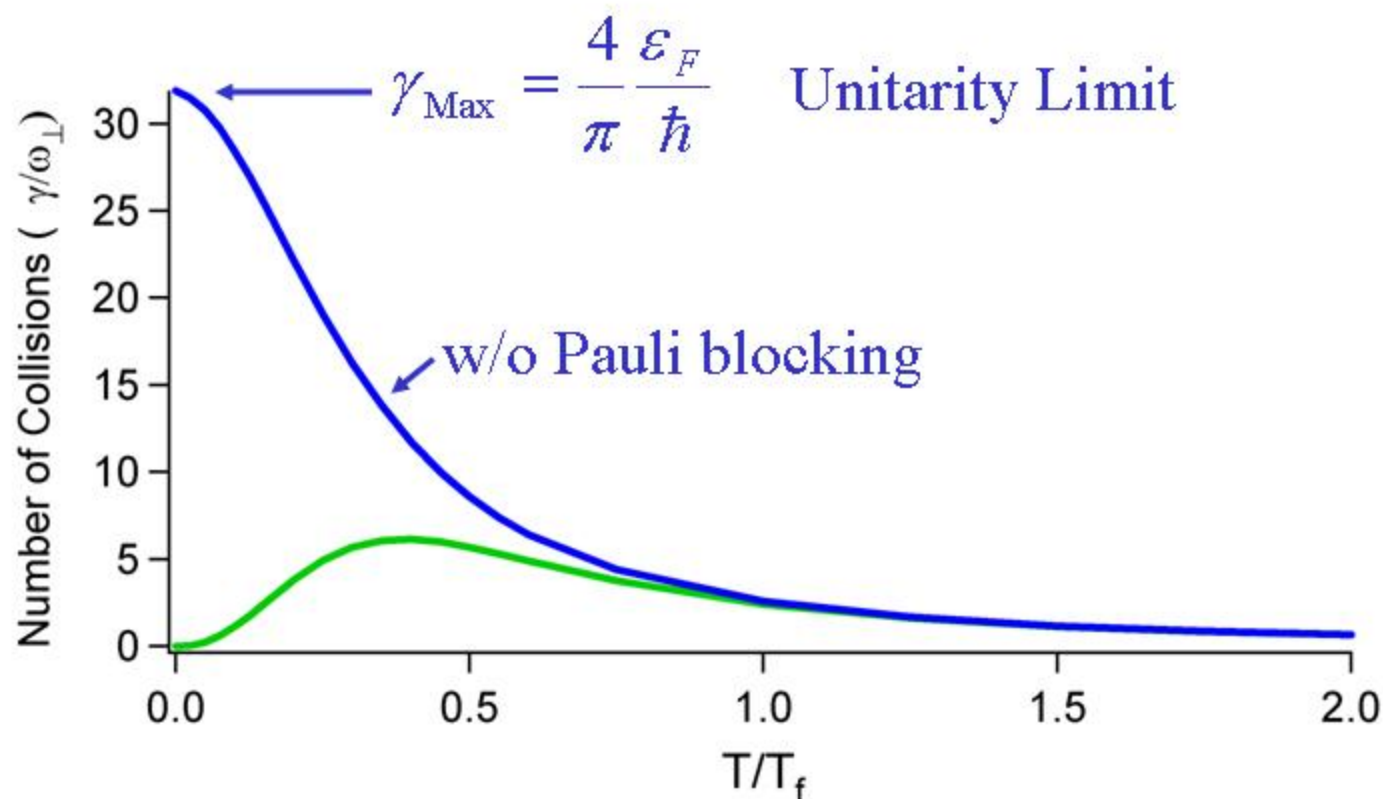
Pauli Blocking in a Harmonic Trap



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Depletion rate **with** and **without** Pauli Blocking:

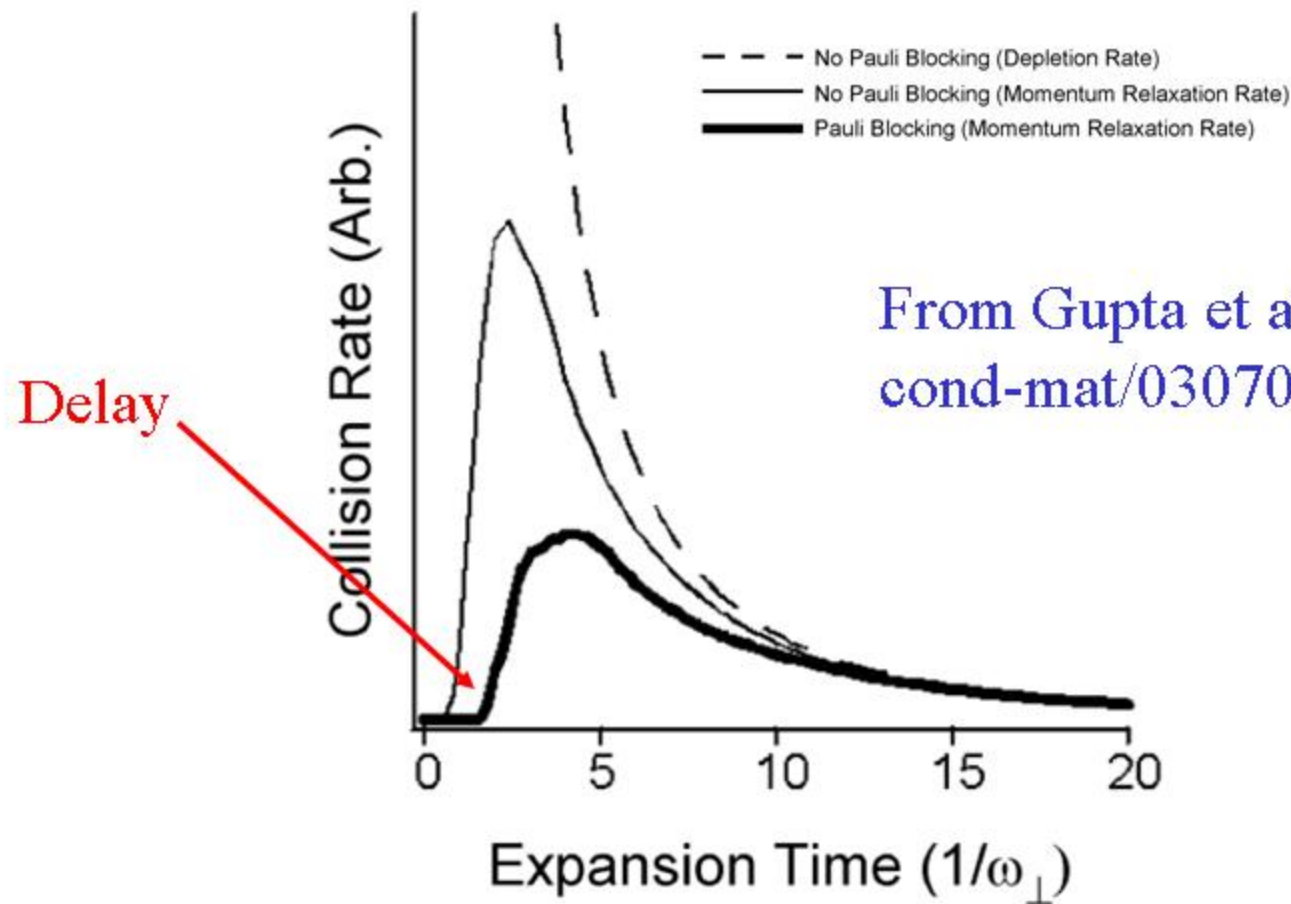


Step-wise Model of Pauli-blocking



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From Gupta et al., arXiv:
[cond-mat/0307088](https://arxiv.org/abs/cond-mat/0307088)

Model: Ballistic expansion for time t_B followed by **No Pauli Blocking**



Relaxation Approximation

Pedri et al., cond-mat/0305624 (2003); Guery-Odelin et al., PRA 60, 4851 (1999)

Ansatz: MB or
Zero temp TF

$$W(\mathbf{x}, \mathbf{p}, t) = \frac{1}{\prod_j (b_j \theta_j^{1/2})} W_0 \left(\tilde{\mathbf{x}}_i, \frac{p_i - M u_i}{\sqrt{\theta_i}} \right) \quad \tilde{\mathbf{x}}_i = \frac{\mathbf{x}_i}{b_i}, \quad u_i = \tilde{\mathbf{x}}_i \dot{b}_i$$

$$0 = \ddot{b}_i + \omega_i^2 b_i [1 + \varepsilon(t)] - \frac{\theta_i}{b_i} \omega_i^2 + \frac{1}{M \langle \tilde{\mathbf{x}}_i^2 \rangle} \left[\frac{1}{b_i} \left\langle \tilde{\mathbf{x}}_i \frac{\partial U_{\text{MF}}(\tilde{\mathbf{x}}, t)}{\partial \tilde{\mathbf{x}}_i} \right\rangle - \frac{\theta_i}{b_i} \left\langle \tilde{\mathbf{x}}_i \frac{\partial U_{\text{MF}}(\tilde{\mathbf{x}}, 0)}{\partial \tilde{\mathbf{x}}_i} \right\rangle \right]$$

Trap Modulation

≈ 0

For hydrodynamic limit
and $U_{\text{MF}}(\mathbf{x}) = \frac{3}{5} \beta \varepsilon_F(\mathbf{x})$

$$\dot{\theta}_i + 2\theta_i \frac{\dot{b}_i}{b_i} = -\frac{1}{\tau} (\theta_i - \bar{\theta}) \quad \bar{\theta} = \frac{1}{3} \sum_j \theta_j$$

No Pauli Blocking:

$$\frac{1}{\tau} = \frac{1}{\prod_j b_j \sqrt{\bar{\theta}}} \frac{8}{25} \left(\frac{4}{\pi} \frac{\varepsilon_F}{\hbar} \right)$$

Momentum relaxation rate

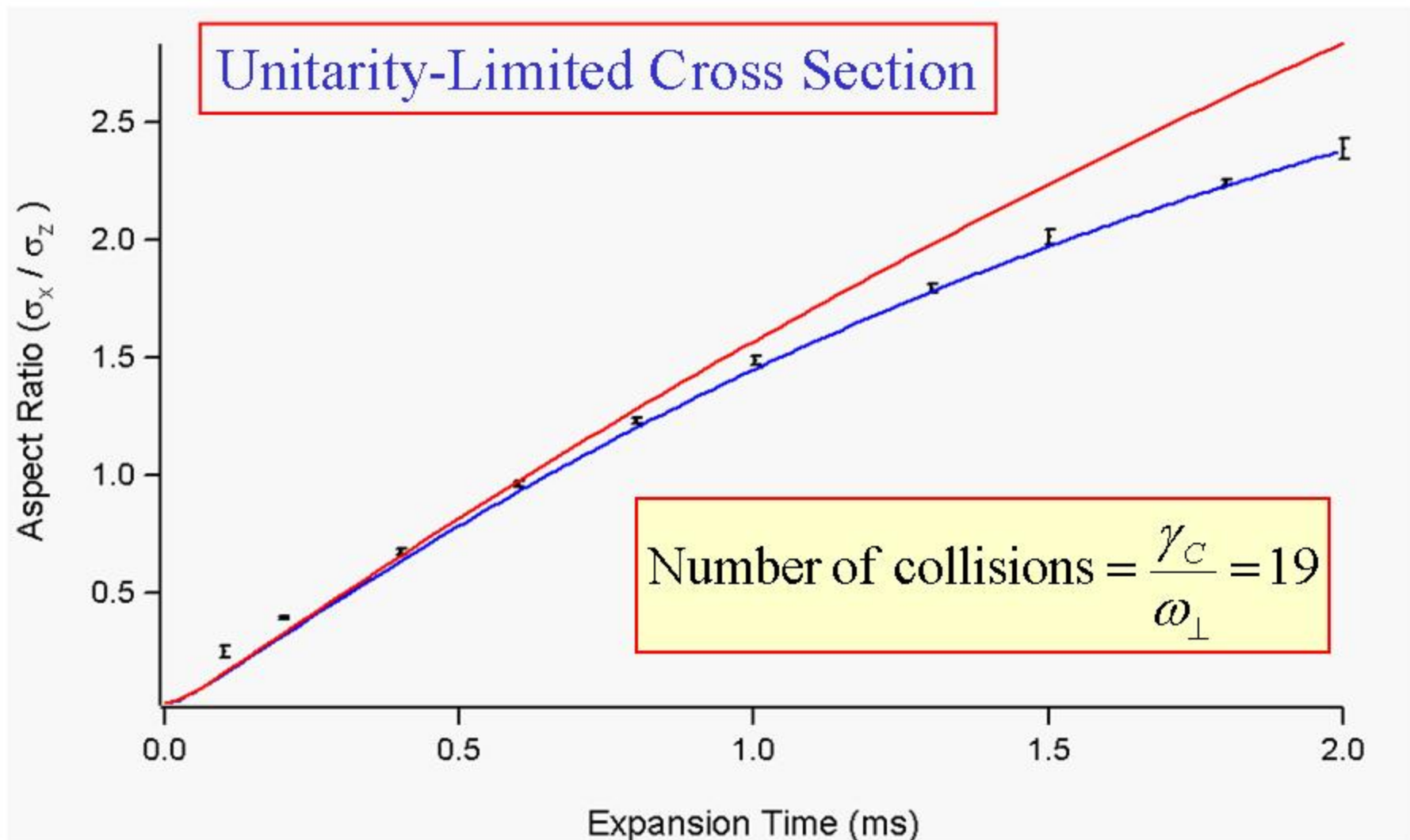
Turn **off** for a time t_B

Collisional Expansion Dynamics



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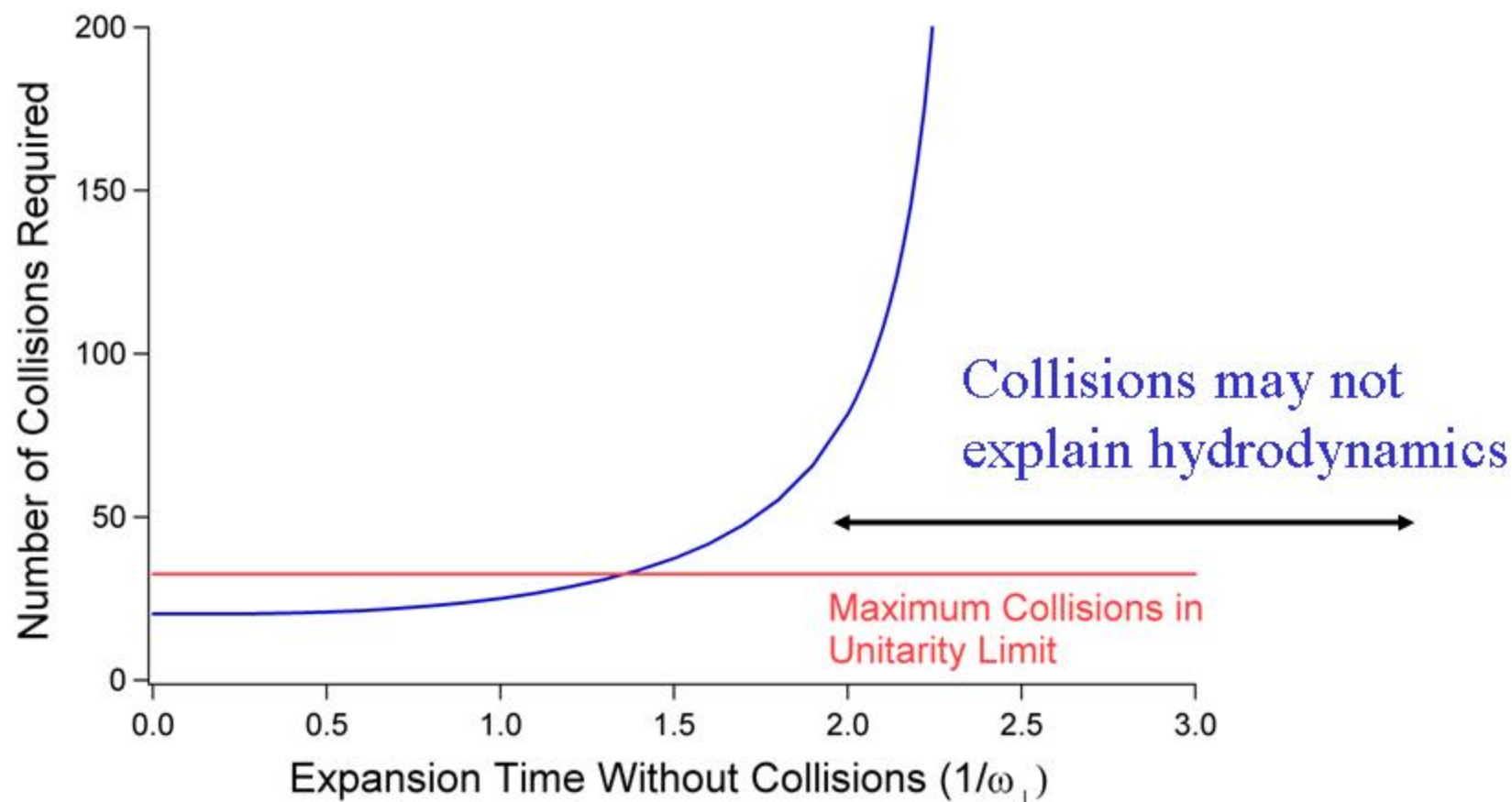
Finite Time to Reinstate Collisions



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$$\text{Number of collisions to achieve our aspect ratio} = \frac{\gamma_c}{\omega_{\perp}}$$



Hydrodynamics in a Trapped Fermi Gas

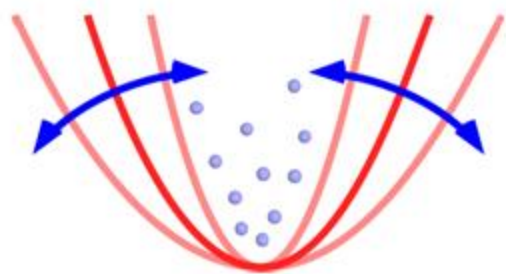


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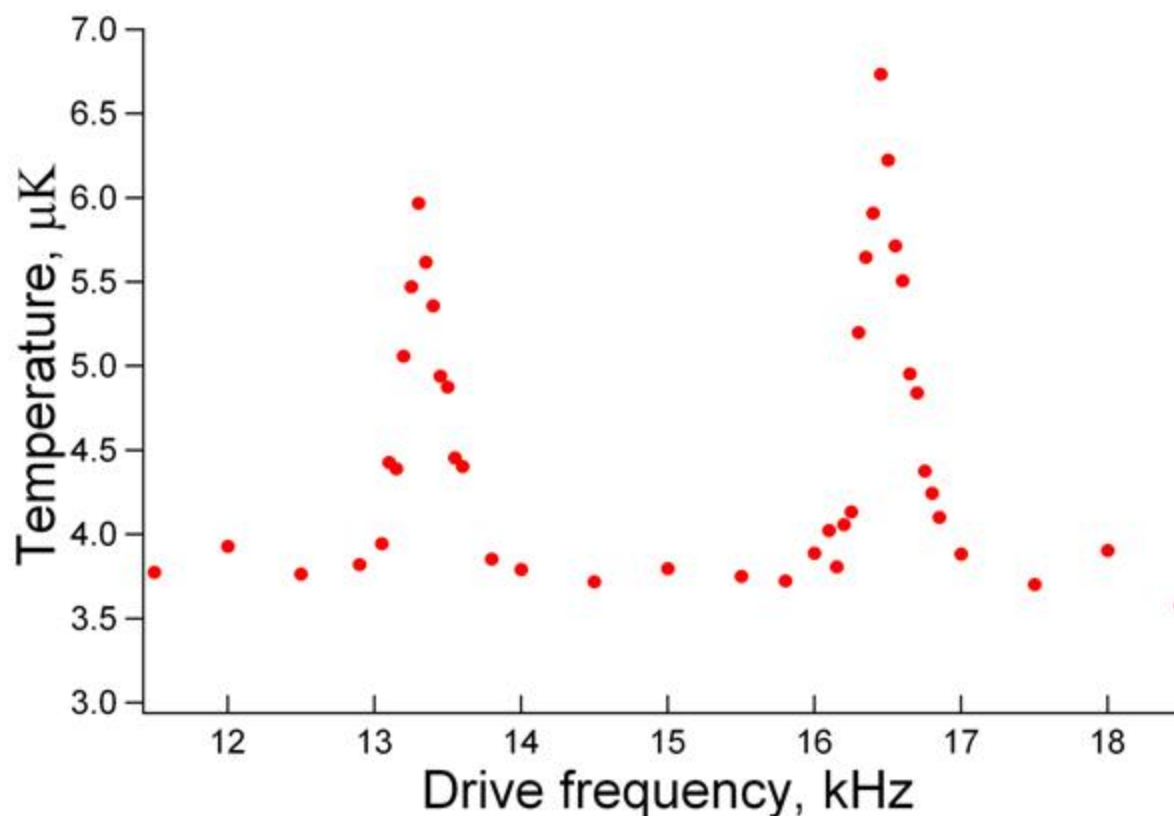
Parametric resonance

$B = 300$ G (few collisions)



Modulate trap depth

Resonance at 2ω



Hydrodynamics in a Trapped Fermi Gas

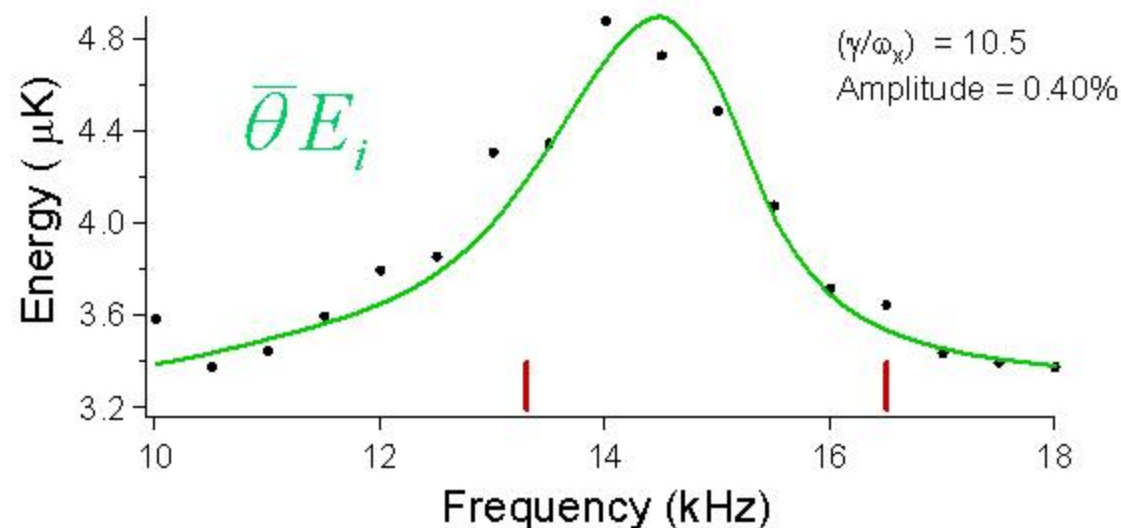
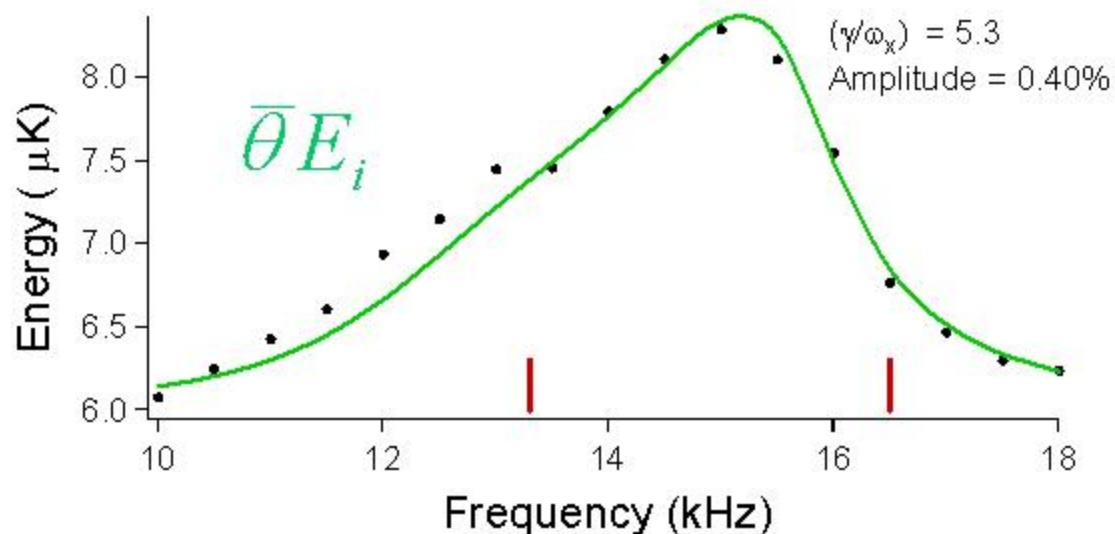


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Atom Cooling and Trapping

$B = 910 \text{ G}$
(Strong Interactions)

$$T = 0.6 T_F$$



Hydrodynamics in a Trapped Fermi Gas

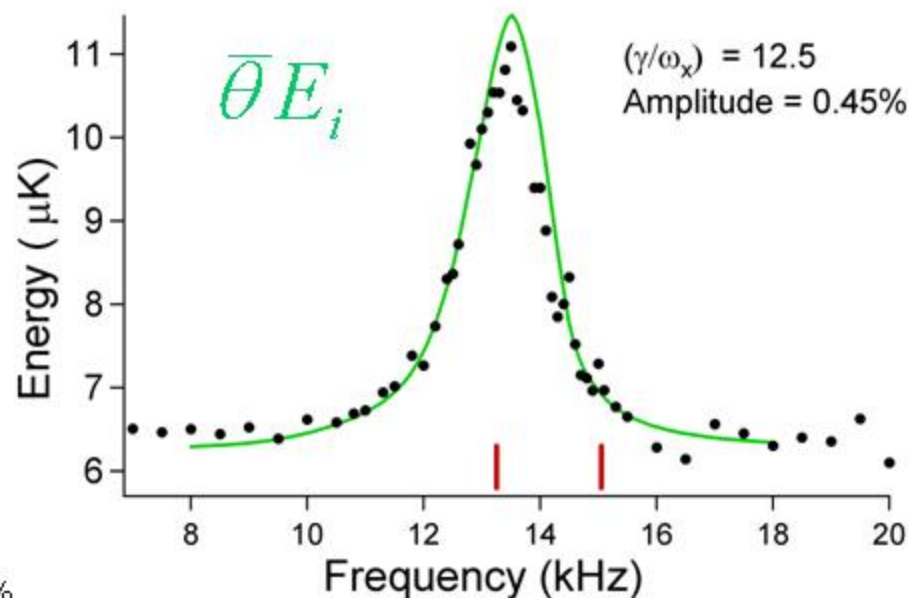
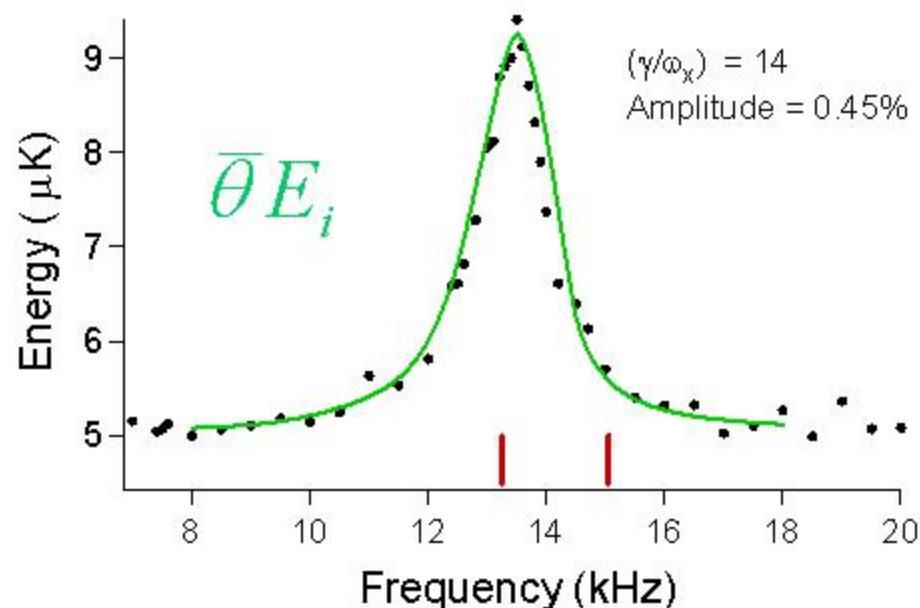


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Atom Cooling and Trapping

$B = 910 \text{ G}$
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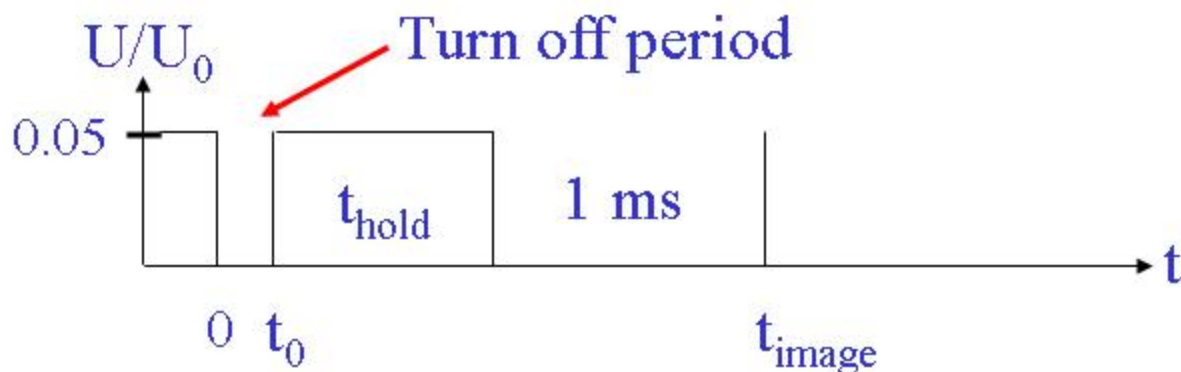
$$T = 0.5 T_F$$

Collective Modes in a Trapped Fermi Gas



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$$t_0 = \frac{1}{16} \frac{2\pi}{\omega} \quad \longrightarrow \quad \Delta E = \frac{1}{2} M \omega^2 \mathbf{v}^2 t_0^2 = \frac{1}{2} (\omega t_0)^2 E = \frac{\pi^2}{128} E$$

$$E = \frac{3}{4} \varepsilon_F \left[1 + \frac{2\pi^2}{3} \left(\frac{T}{T_F} \right)^2 \right] \quad \frac{T_i}{T_F} \cong 0.1 - 0.15 \quad \longrightarrow \quad \boxed{\frac{\Delta T}{T_F} \cong 0.05}$$

Collective Modes in a Trapped Fermi Gas

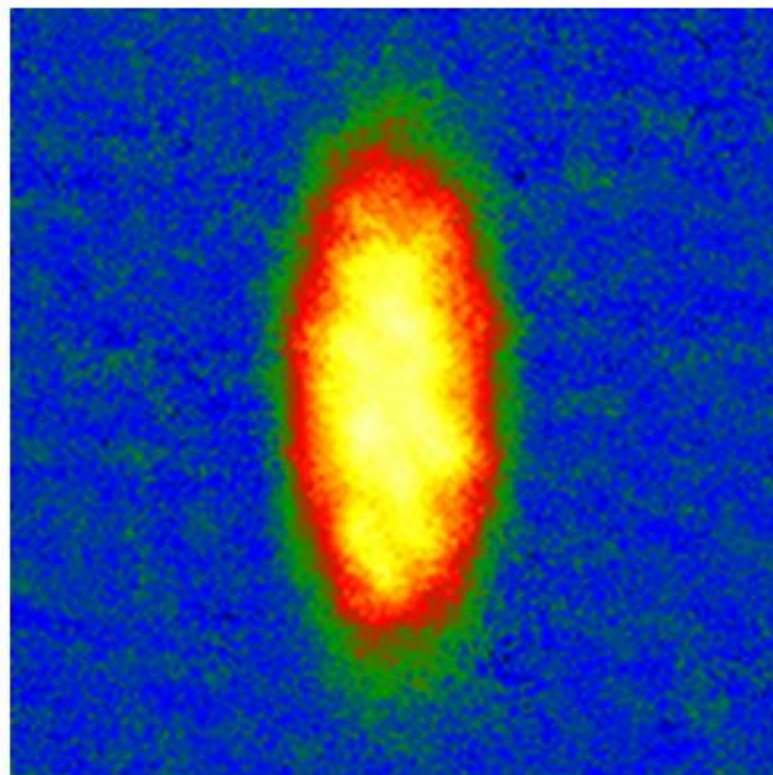


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Atom Cooling and Trapping

$$B = 870 \text{ G}$$

$$\frac{T_i}{T_F} = 0.1 - 0.15$$



Collective Modes in a Trapped Fermi Gas



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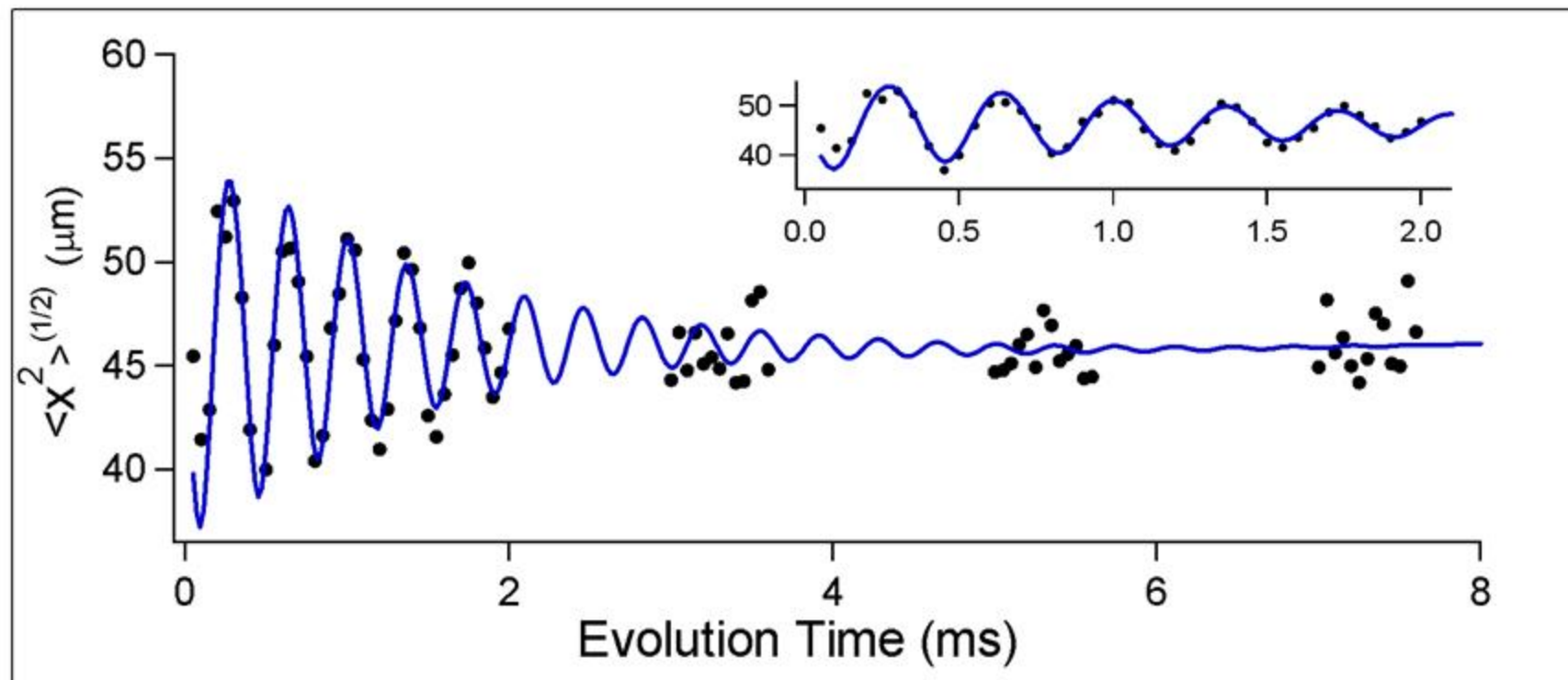
Hydrodynamic:

Fit Parameters

$$\begin{aligned} 1/(\omega_x \tau_0) &= 4.8 & (\langle x^2 \rangle_0)^{(1/2)} &= 4.1 \text{ } \mu\text{m} \\ \gamma/\omega_x &= 14 & \nu_x &= 1520 \text{ Hz} \end{aligned}$$

$$B = 870 \text{ G}$$

$$\frac{T_i}{T_F} = 0.3$$



Collective Modes in a Trapped Fermi Gas



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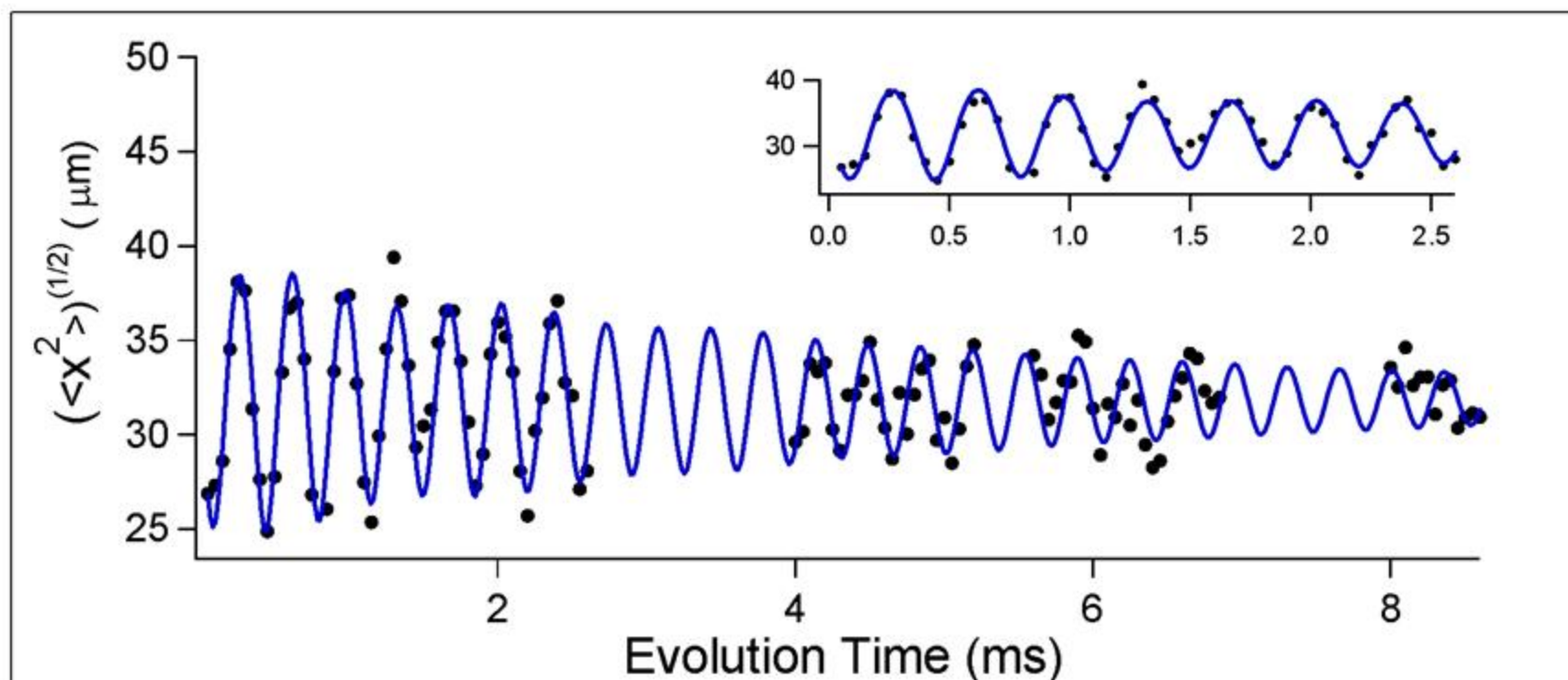
Hydrodynamic:

Fit Parameters

$$\begin{aligned} 1/(\omega_x \tau_0) &= 20 & (\langle x^2 \rangle_0)^{(1/2)} &= 2.6 \text{ } \mu\text{m} \\ \gamma/\omega_x &= 63 & \nu_x &= 1600 \text{ Hz} \end{aligned}$$

$$B = 870 \text{ G}$$

$$\frac{T_i}{T_F} = 0.1 - 0.15$$



Collective Mode Frequencies



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Measured Parametric Resonance Frequencies:

$$\omega_x = 2\pi \times 1610 \text{ Hz}$$

$$\omega_y = 2\pi \times 1510 \text{ Hz}$$

$$\omega_{\perp} = \sqrt{\omega_x \omega_y} = 2\pi \times 1550 \text{ Hz}$$

Predicted Frequency for **Hydrodynamic Fermi** Gas:

$$\omega_{\text{Fermi}} = \sqrt{\frac{10}{3}} \omega_{\perp} = 2\pi \times 2.83 \text{ kHz}$$

Measured Oscillation Frequency for Sinusoidal Fit:

$$\omega_{\text{Meas}} = 2\pi \times 2.84 \text{ kHz}$$

- All-Optical Production of Degenerate Fermi Gas
 - Efficient evaporation near Feshbach resonance
 - Very low T/T_F

Hydrodynamics of a Strongly-Interacting Fermi Gas

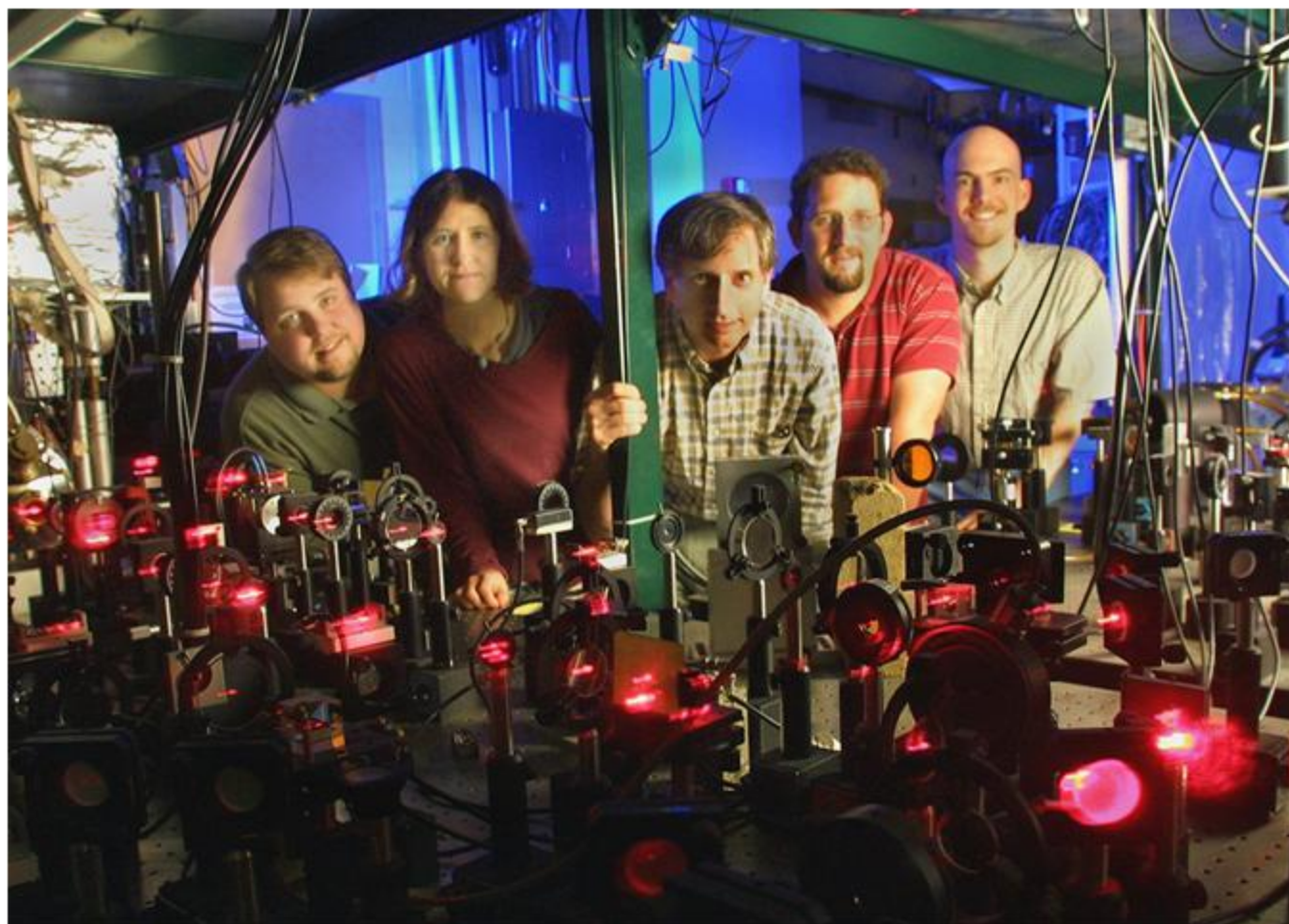
- Observation of Anisotropic Expansion
 - Relaxation approximation model – no Pauli blocking
 - Step-wise model: ballistic drift time t_B + no Pauli blocking
 - For low T , $t_B > 1.5 / \omega_{\perp}$, collisions may not explain hydrodynamics
- Trapped Atom Hydrodynamics
 - Collisionally-damped hydrodynamic modes at high T
 - Weakly-damped hydrodynamic modes at low T
 - Hydrodynamic Fermi gas at $T = 0.1$: Fermi Superfluid?

The Team



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Mike Gehm
Staci Hemmer
John Thomas
Ken O'Hara
Stephen Granade

Joe Kinast
Andrey Turlapov