

Spin polarons and molecules in strongly interacting atomic Fermi gases

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work done in collaboration with:

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Pairing and superfluidity in Fermi systems

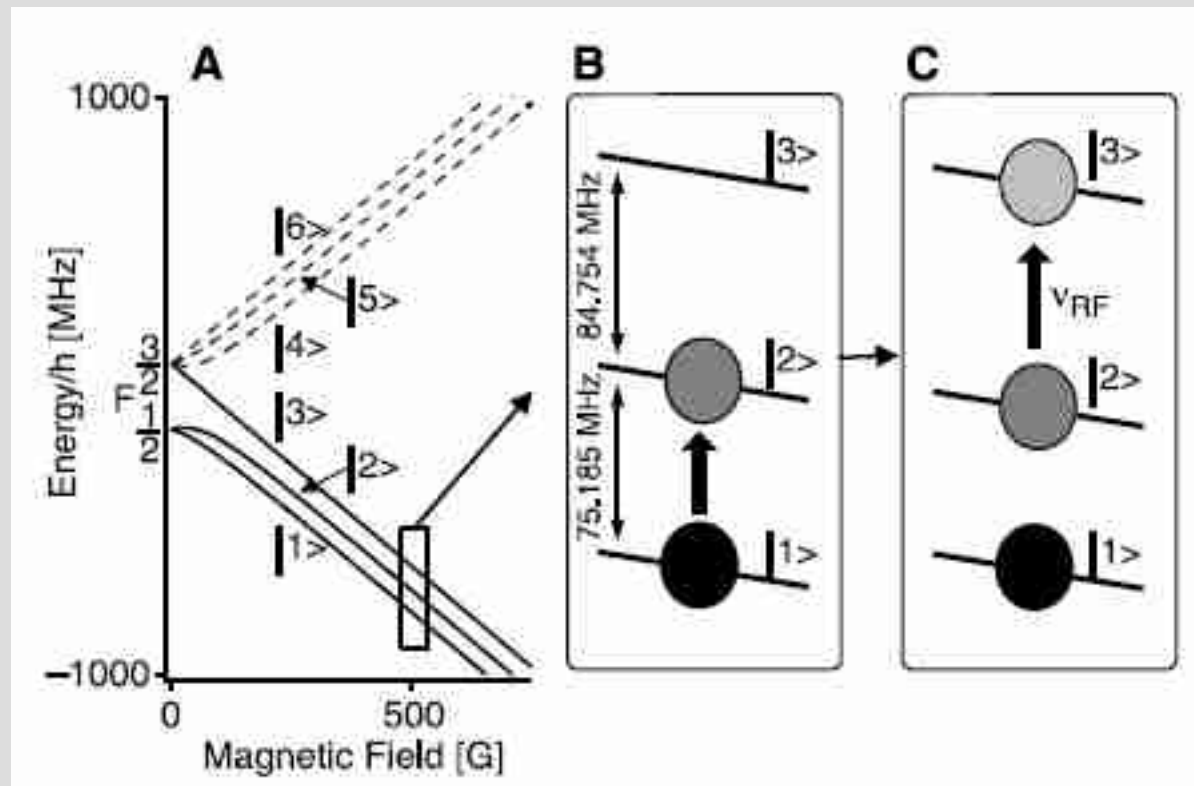
- Strongly-correlated Fermi gases:
 - ♦ neutron stars (huge and hot)
 - ♦ quark-gluon plasma and high- T_c superconductors
 - ♦ atomic gases (tiny and ultracold)
- In BCS theory pairs condense as they form, but when interactions are strong $T_{\text{pairing}} > T_{\text{superfl.}}$
- **Rf measurements**: a tool to probe interactions, onset of pair correlations and presence of molecules.

C. Regal et al. (Nature 2003), C. Chin et al., (Science 2004).

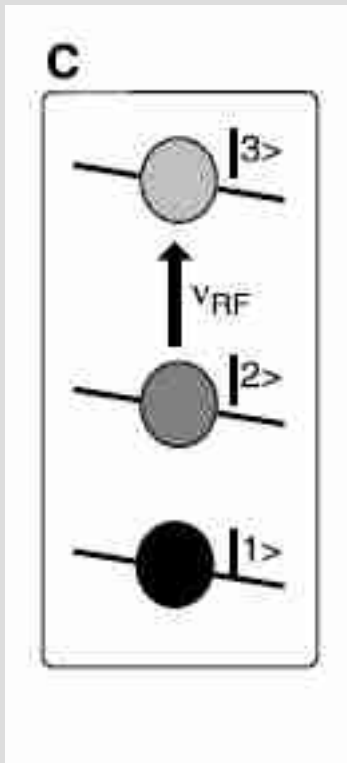
By the end of the talk I wish you will know more about:

- Basics of RF spectroscopy
- Our model for the normal state ($T > T_F$)
- What can we learn from the spectral function?
- Spectra of balanced and imbalanced systems
- How many pairs in the BEC-BCS crossover?
- An impurity in a Fermi sea: polaron vs. molecule
- (... + topological modes and Dirac fermions in BF mixture)

Level structure in ${}^6\text{Li}$



Radio-frequency spectroscopy



$T \gg T_F \rightarrow$

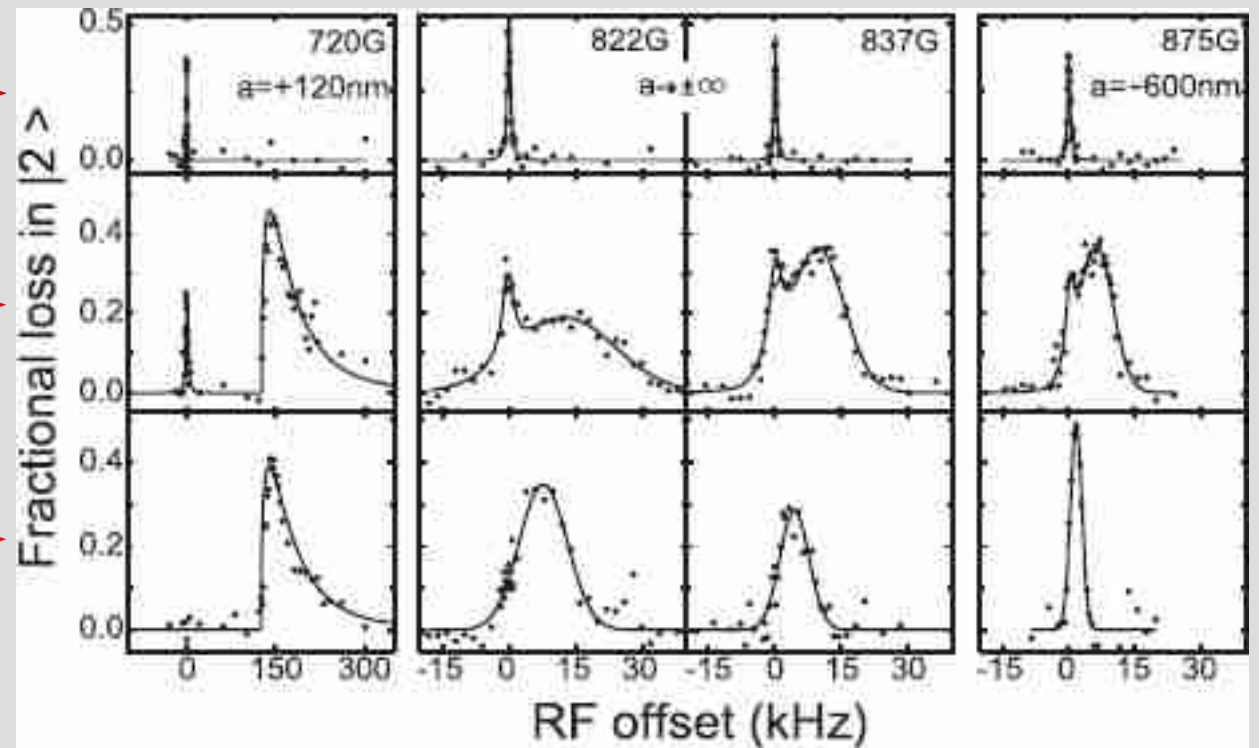
$T \sim 0.5 T_F \rightarrow$

$T \ll T_F \rightarrow$

BEC

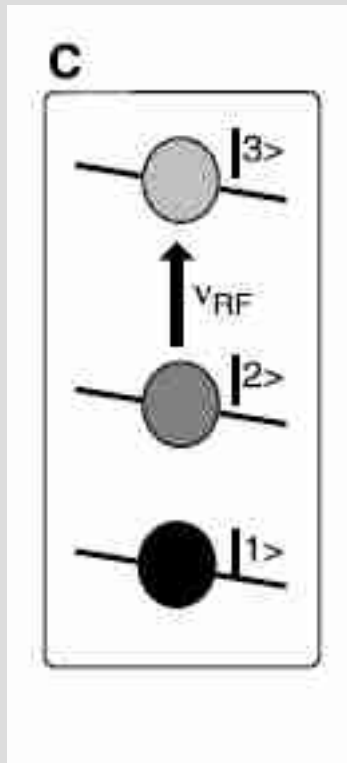
Unitarity

BCS



C. Chin et al., (Science 2004).

Linear response theory



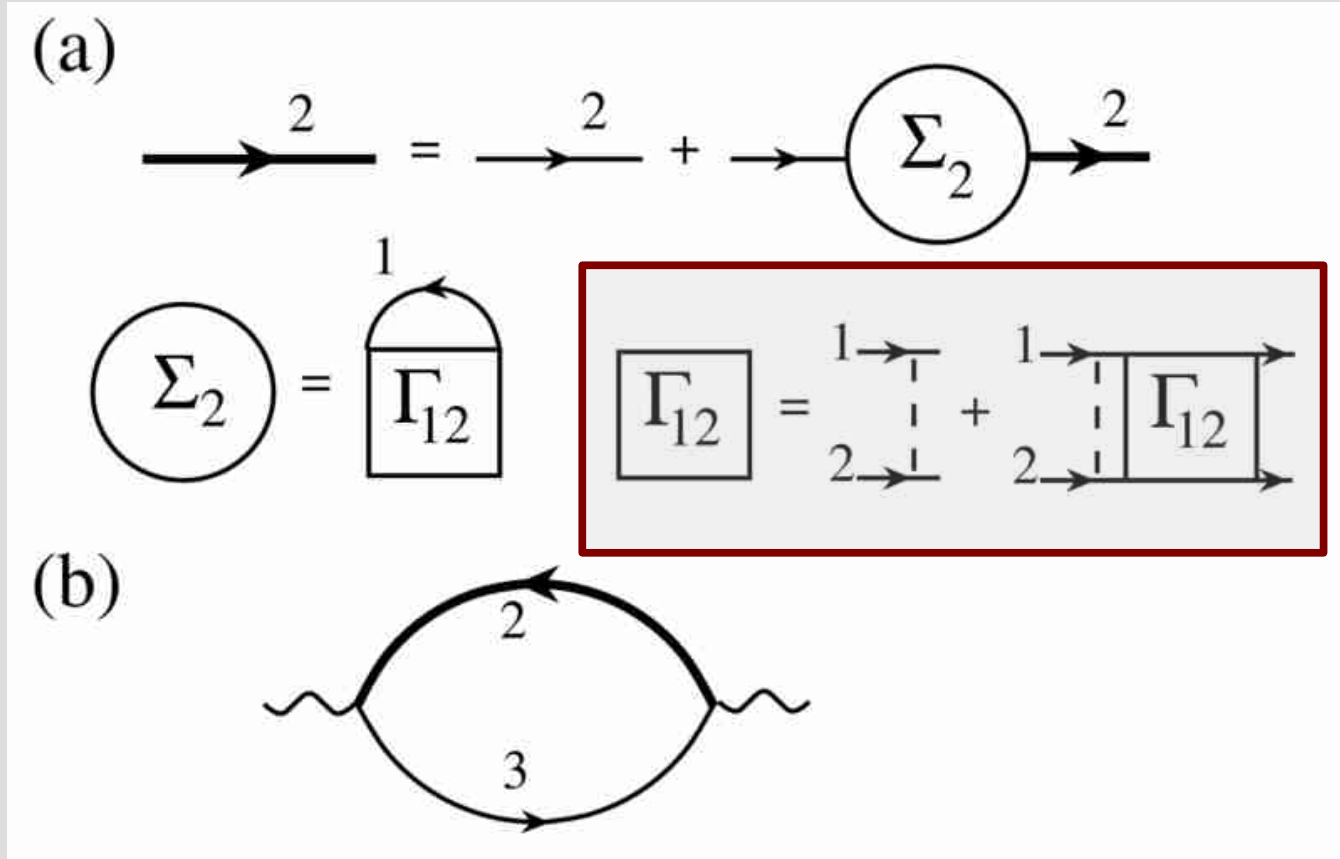
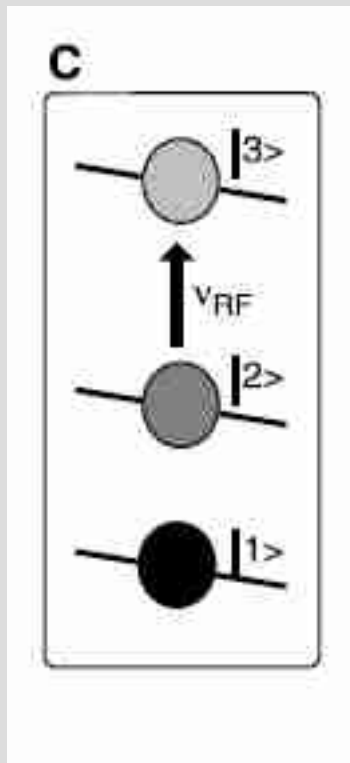
$$H_{rf} = \frac{\Omega}{2} \int d\mathbf{r} \left[e^{-i\omega t} \psi_3^\dagger(\mathbf{r}, t) \psi_2(\mathbf{r}, t) + \text{h.c.} \right]$$

$$\text{Transition rate: } R(\omega) \propto - \int d\mathbf{r} d\mathbf{r}' \text{Im} \mathcal{D}(\mathbf{r}, \mathbf{r}', \omega)$$

$$\mathcal{D}(\mathbf{r}, \mathbf{r}', \omega) = \text{F.T.} \left\{ -i\theta(t - t') \langle [\psi_3^\dagger(\mathbf{r}, t) \psi_2(\mathbf{r}, t), \psi_2^\dagger(\mathbf{r}', t') \psi_3(\mathbf{r}', t')]] \rangle \right\}.$$

{...}: retarded 2-3 flip correlation function

T-matrix and self-energy in the ladder approx.



Spectral function $A_2(k, \omega) = -2\text{Im}[G_2(k, \omega)]$

(we take into account the strong 1-2 interaction, while we neglect the weaker 1-3)

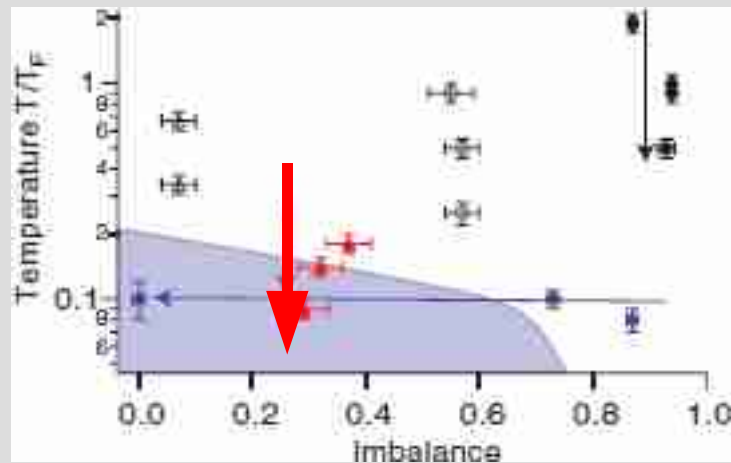
Pairing Without Superfluidity: The Ground State of an Imbalanced Fermi Mixture

C. H. Schunck,* Y. Shin, A. Schirotzek, M. W. Zwierlein,† W. Ketterle

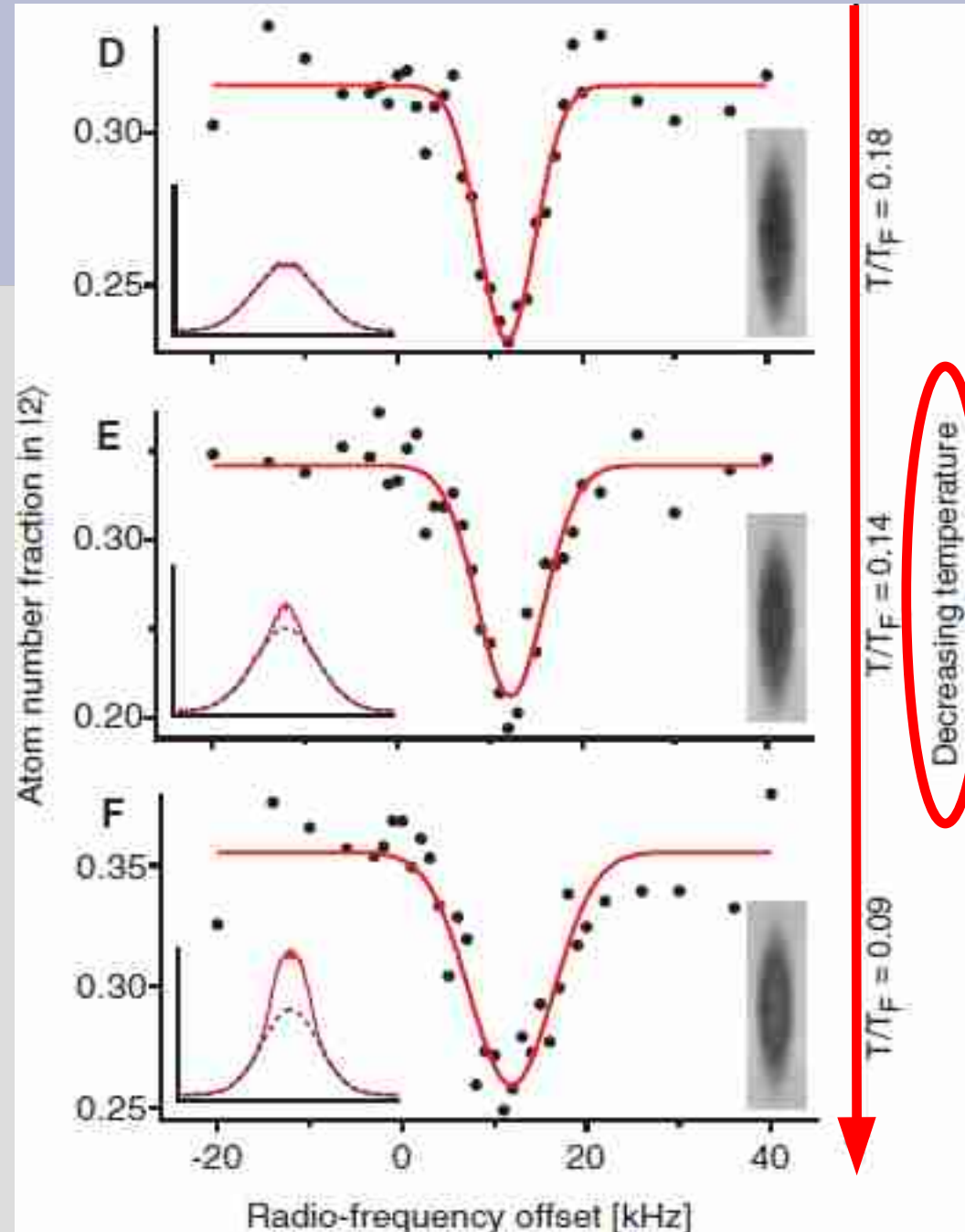
Science (2007)

RF spectra do not reveal
the normal-SF transition

(while density profiles do)



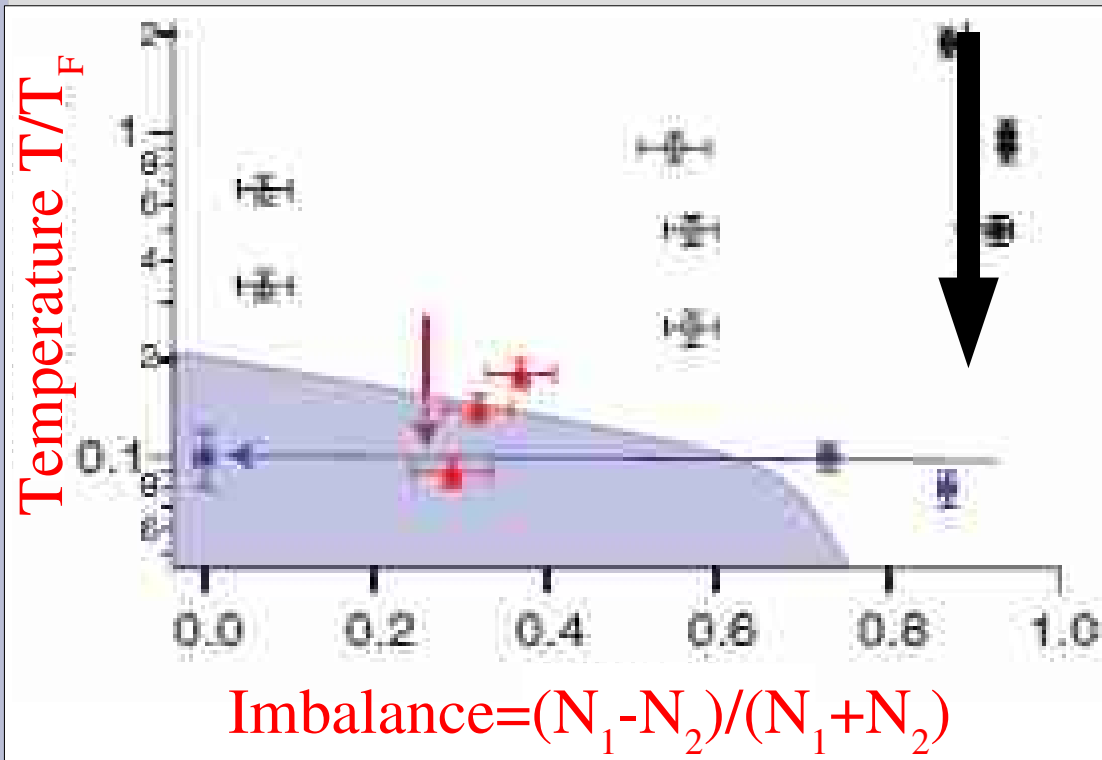
$$\text{Imbalance} = (N_1 - N_2) / (N_1 + N_2)$$



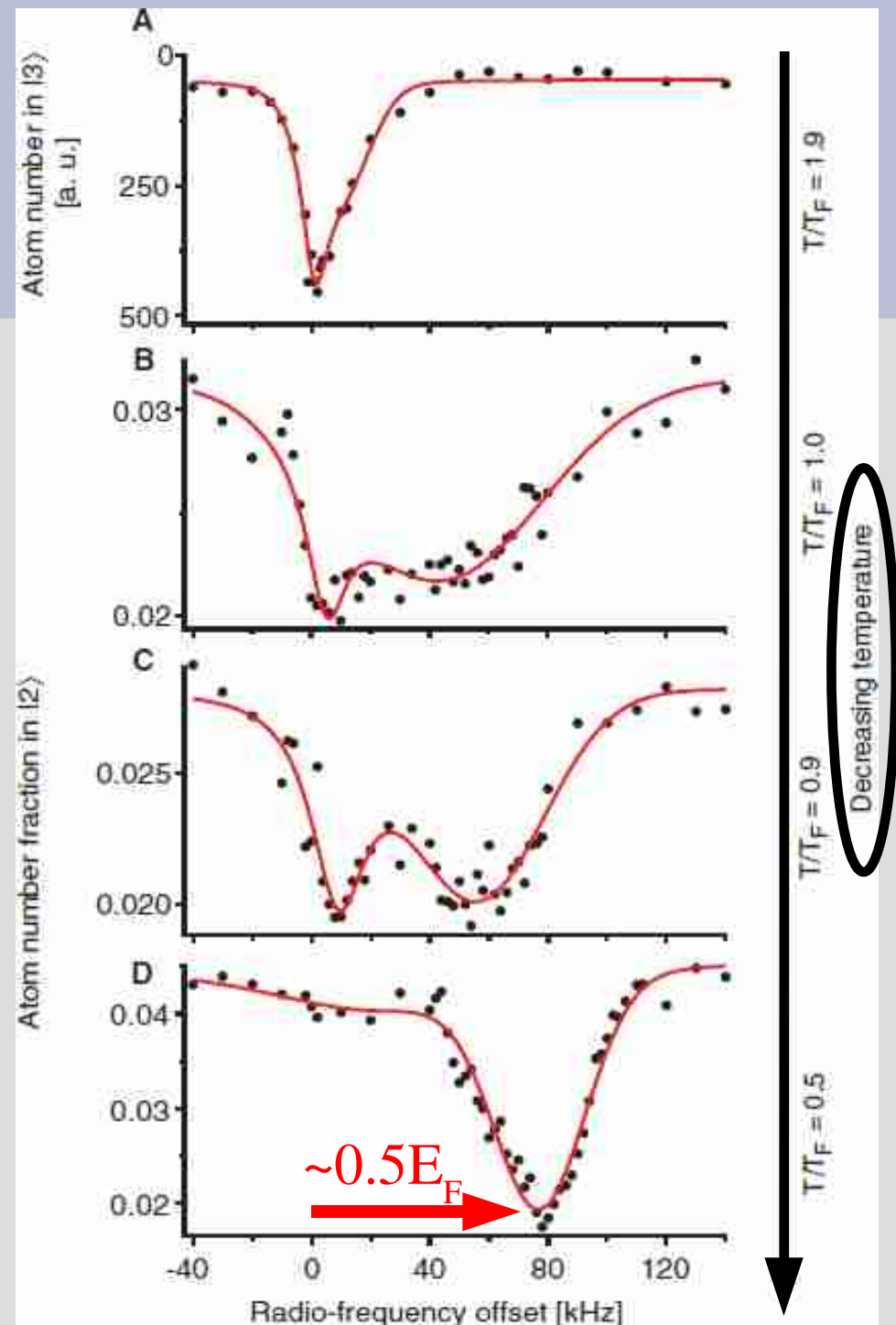
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Science (2007)



Large shift in the peak position
signals strong interactions in the gas



Spectral function and spectrum (no trap, unitarity)

$$\text{Im}D(\omega) = -\mathcal{V} \int \frac{d\mathbf{k}}{(2\pi)^3} A_2(k, \xi_{2k} - \omega) f(\xi_{2k} - \omega)$$

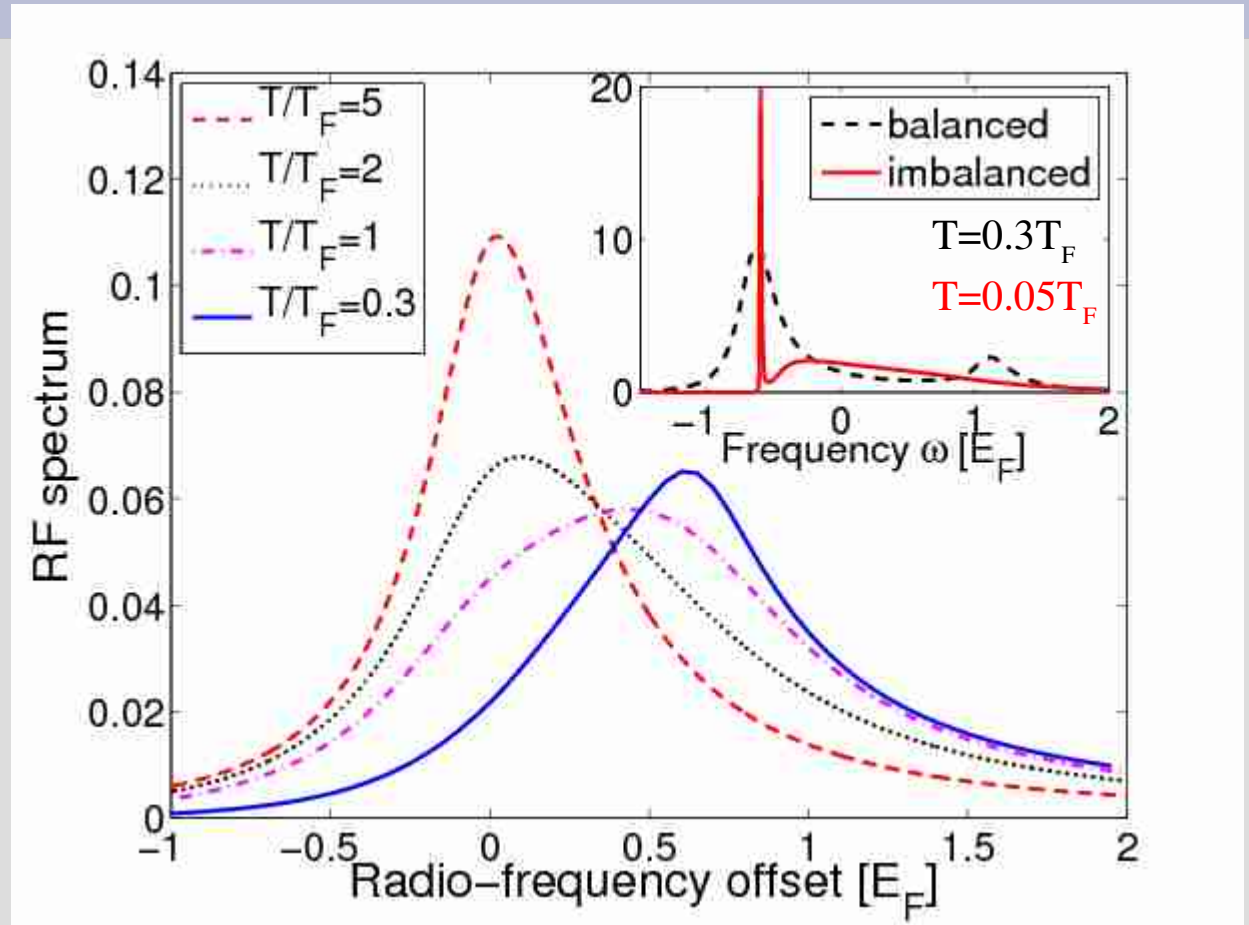
$$\xi_{2k} = k^2/2m - \mu_2$$

$$\text{Spectral function: } A_2(k, \omega) = -2\text{Im}G_2(k, \omega)$$

Pairing
without
superfluidity
(*pseudogap* regime)

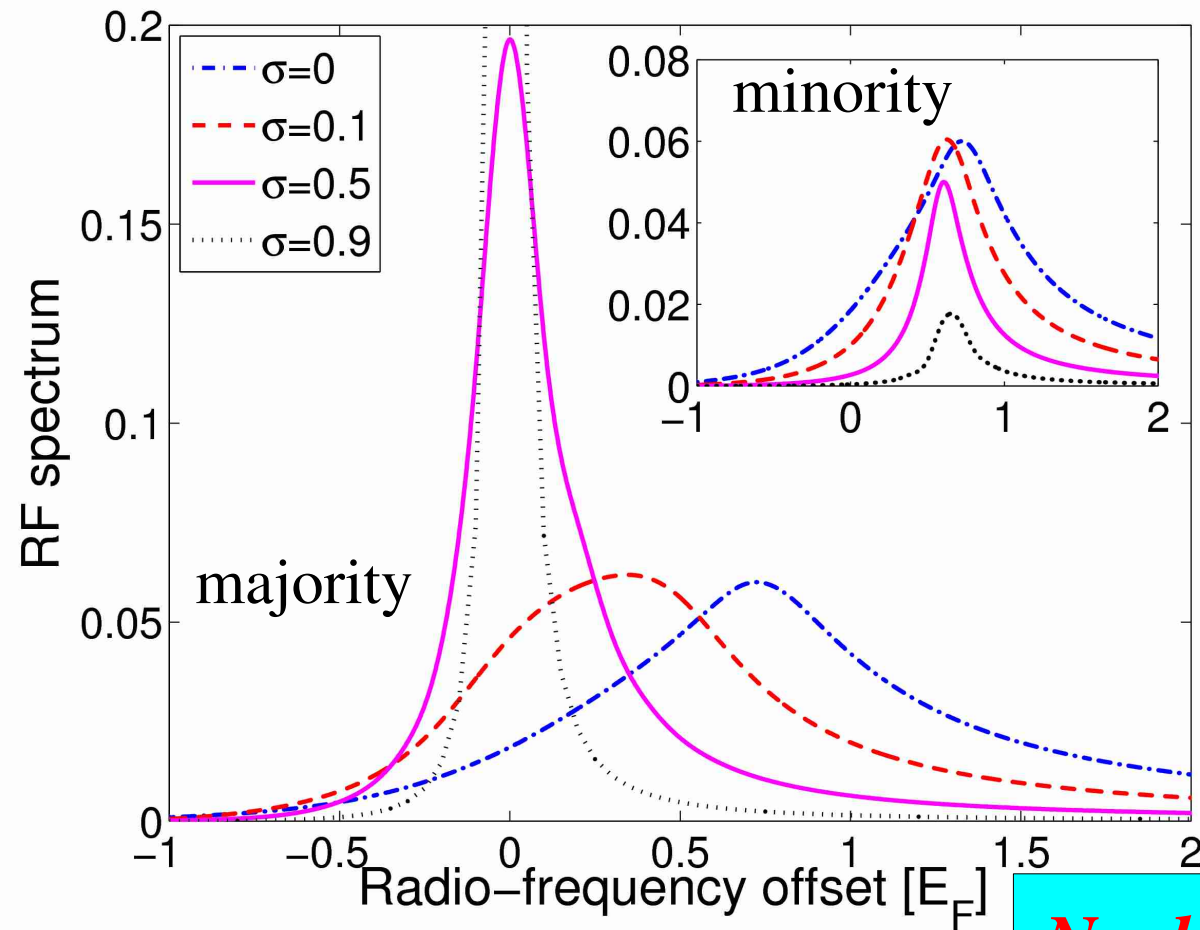
Inset: spectral function $A_2(k=0, \omega)$.

Main: spectrum for $n_1=n_2$.



***No double peaks in the spectrum,
even in the pseudogap regime!***

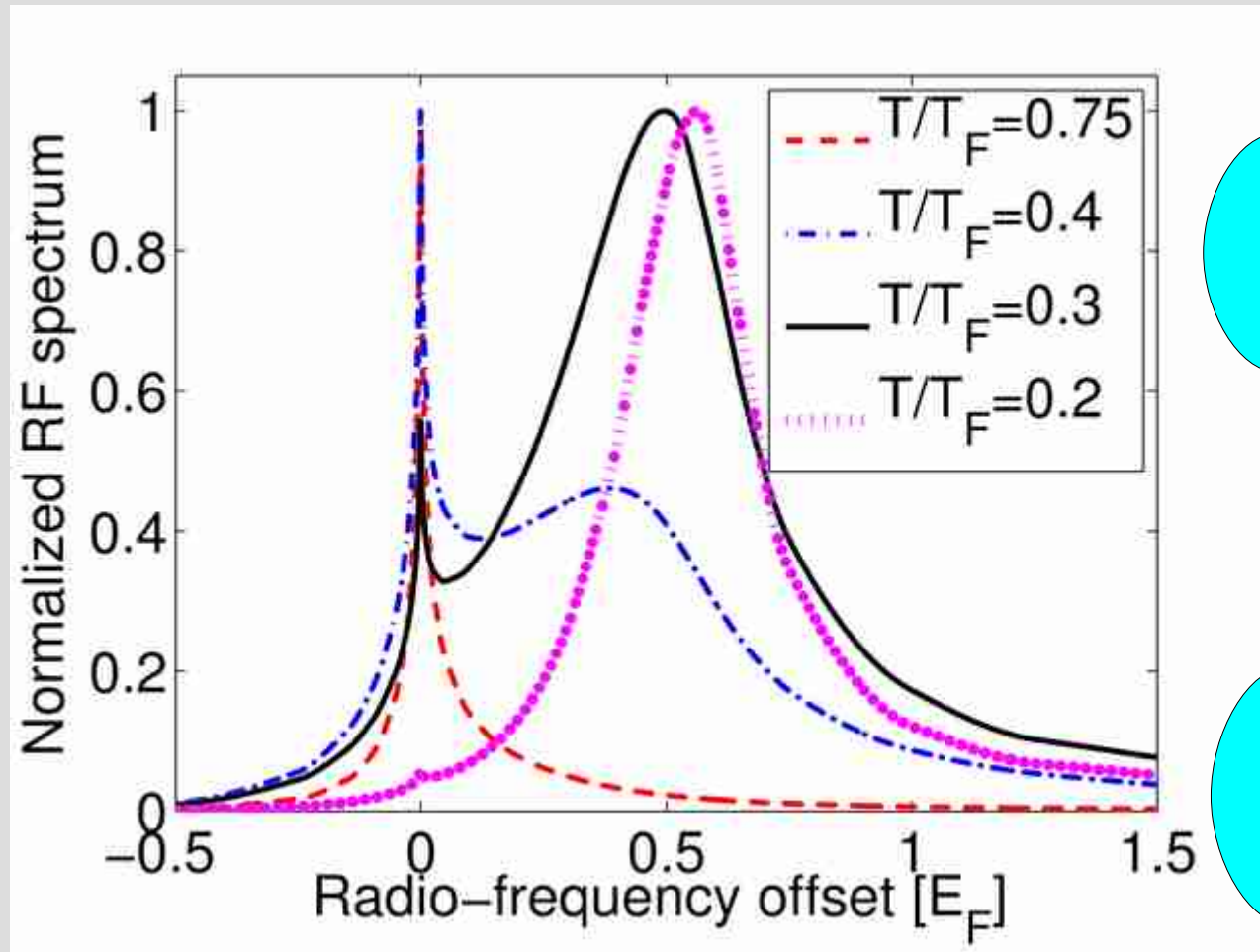
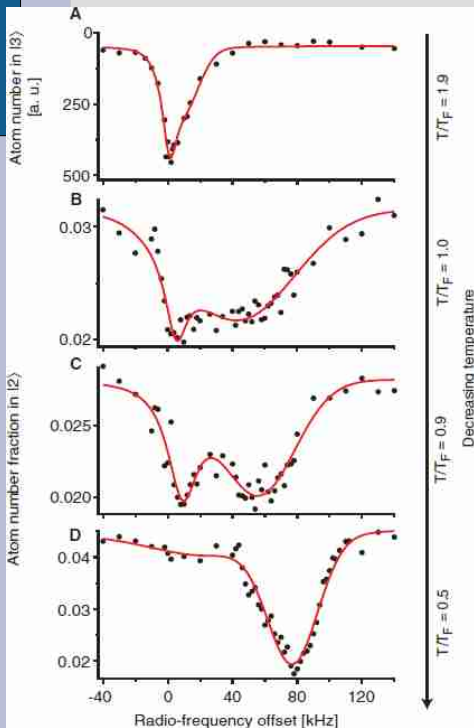
Spectrum of an imbalanced gas at unitarity



minority peak position
independent
of imbalance

*No double peaks in the spectrum,
even in the pseudogap regime!*

Minority spectrum of an imbalanced **trapped** gas at unitarity

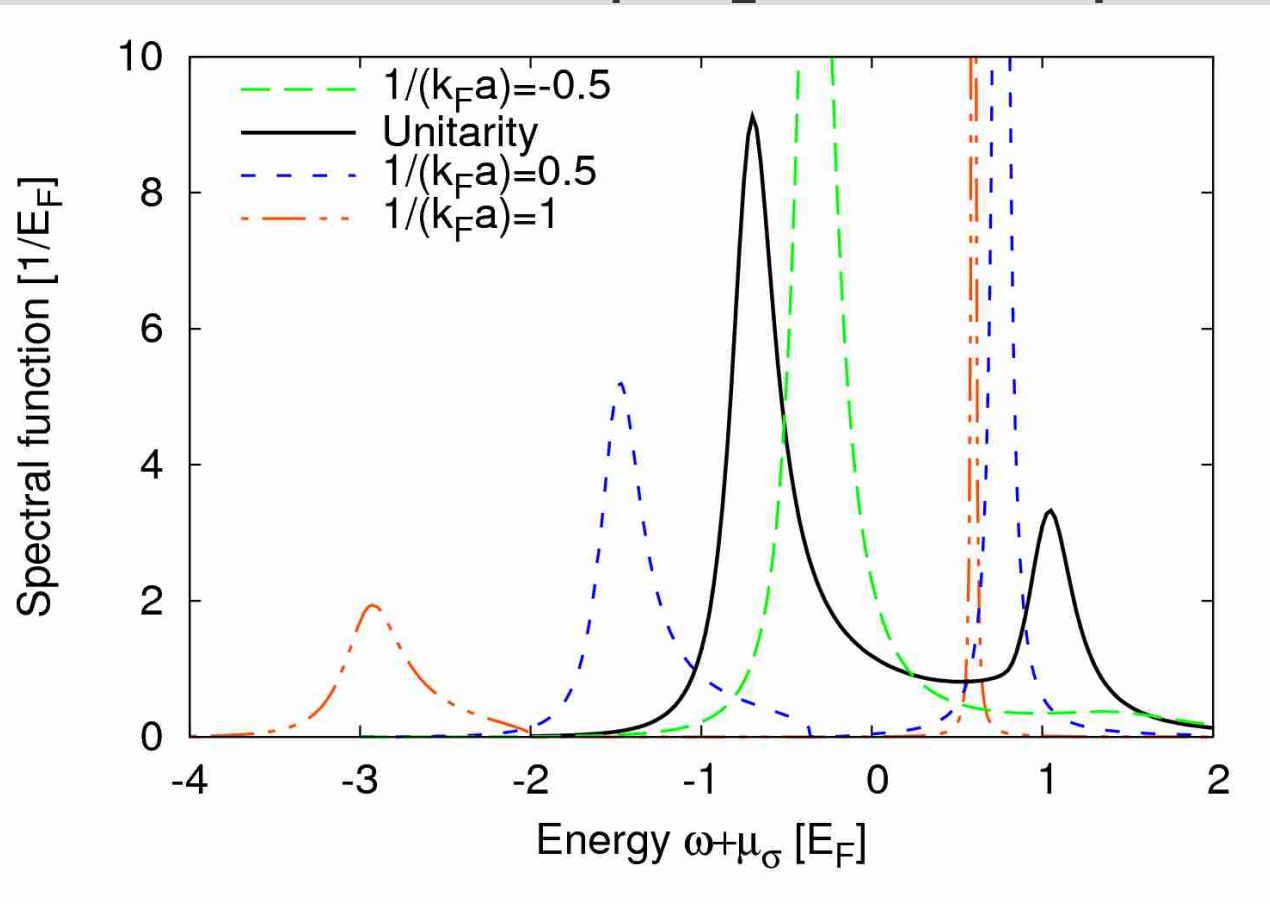


no fitting parameters

peak position independent of imbalance

$$N_2/N_1 = 5\%$$

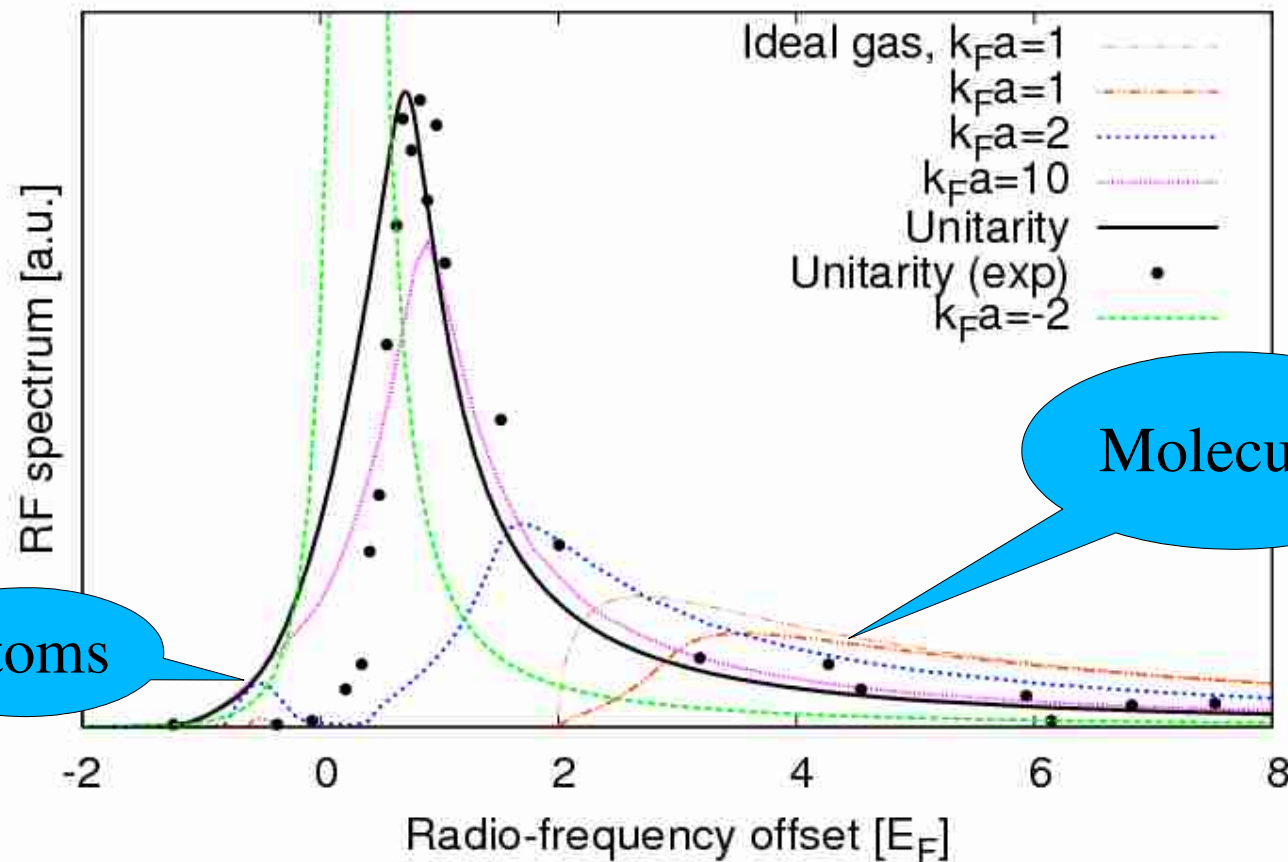
Spectral function $A_2(k=0, \omega)$ across the crossover (no trap, $n_1=n_2$, $T=0.3T_F$)



a real gap in the spectral function opens for $1/(k_F a) > 0.5$

Spectrum across the crossover

(no trap, $n_1=n_2$, $T=0.3T_F$)



1) Stable molecules appear when a gap opens in A_2

2) Atoms coexist with molecules for $T < E_b$

Molecules

Atoms

$$\text{Binding energy: } \frac{E_b}{E_F} \sim \frac{2}{(k_F a)^2}$$

Exp: Schunck et al., Nature 2008

(thin line: ideal gas model, Yukawa molecular function)

Number of “pairs/molecules” δn_σ in the ladder (NSR) approx.

Thermodynamic potential: $\Omega = \Omega_0 + \Delta\Omega$

NSR approx:
$$\Delta\Omega = \int \frac{d^3k}{(2\pi)^3} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi i} \frac{\ln(1 - T\Pi^+) - \ln(1 - T\Pi^-)}{e^{\beta\omega} - 1}$$

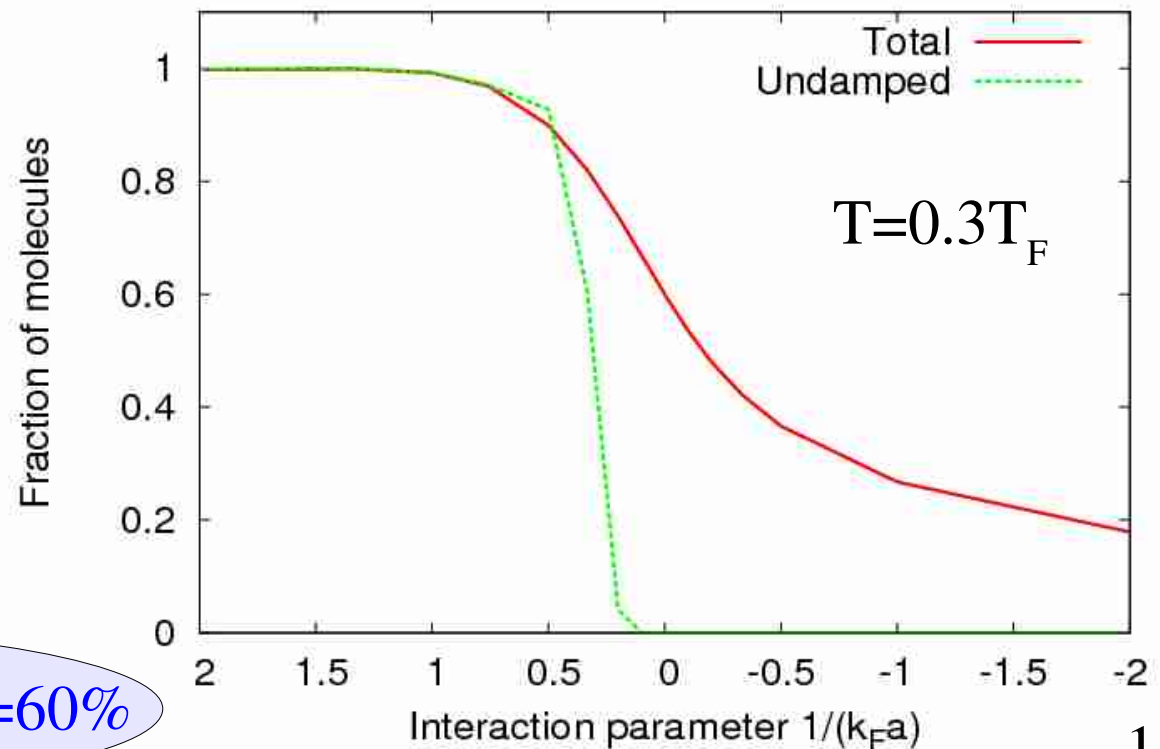
$$n_\sigma = -\partial_{\mu_\sigma} \Omega =$$

$$= n_\sigma^0 + \delta n_\sigma$$

$$n_\sigma^0 = \int \frac{d^3k}{(2\pi)^3} [e^{\beta(k^2/2m - \mu)} + 1]^{-1}$$

$$\delta n_\sigma = -\partial_{\mu_\sigma} \delta\Omega$$

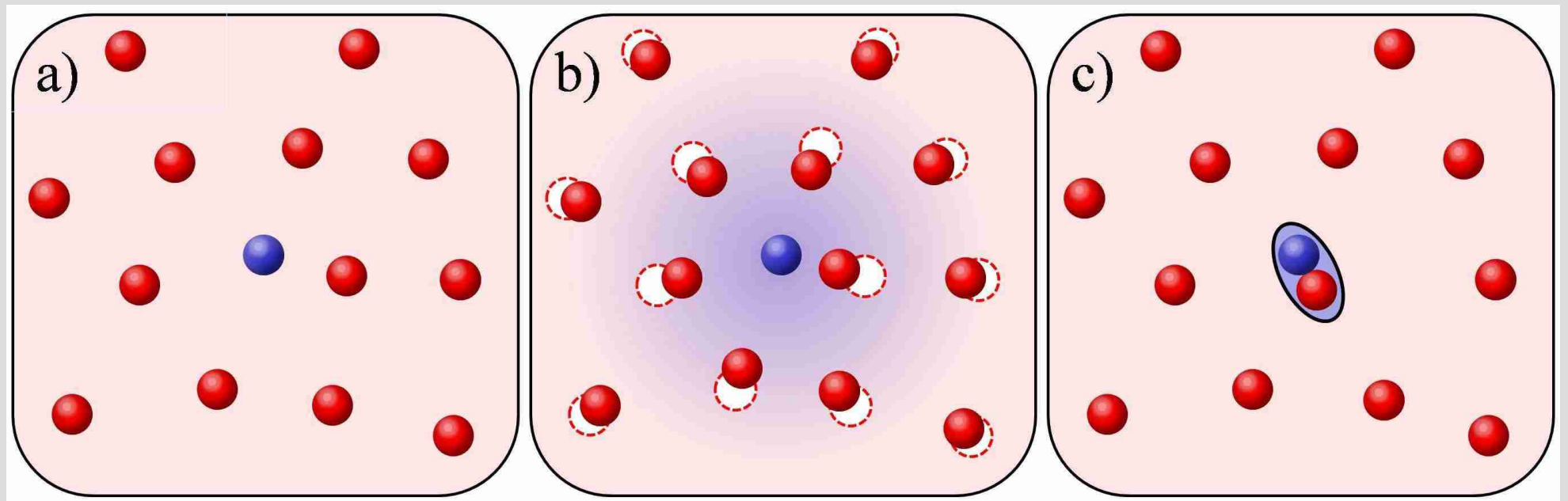
($\sigma = 1, 2$)



At $T=T_c$ and unitarity, $\delta n_\sigma = 60\%$

An impurity in a Fermi sea (the MIT experiment)

increasing attractive interaction

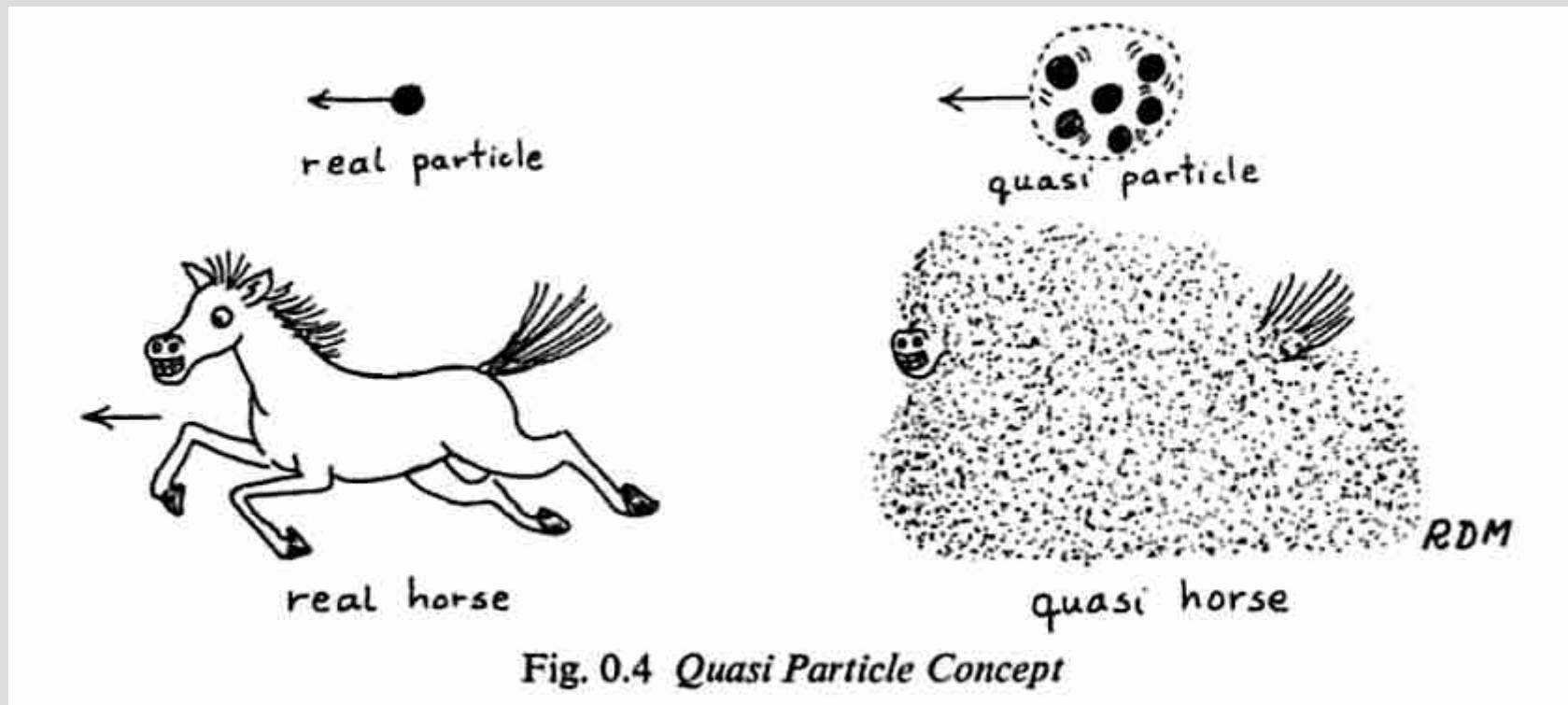


a free particle

a polaron

a molecule

The impurity problem (in Mattuck's imagination)



The impurity @ FerMix

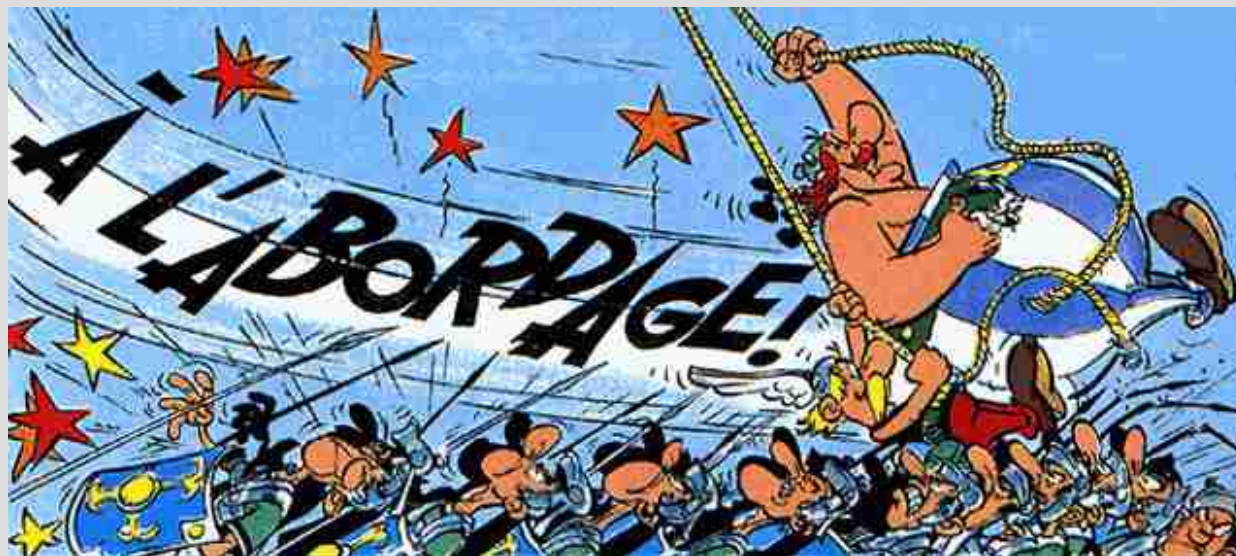


the free particle

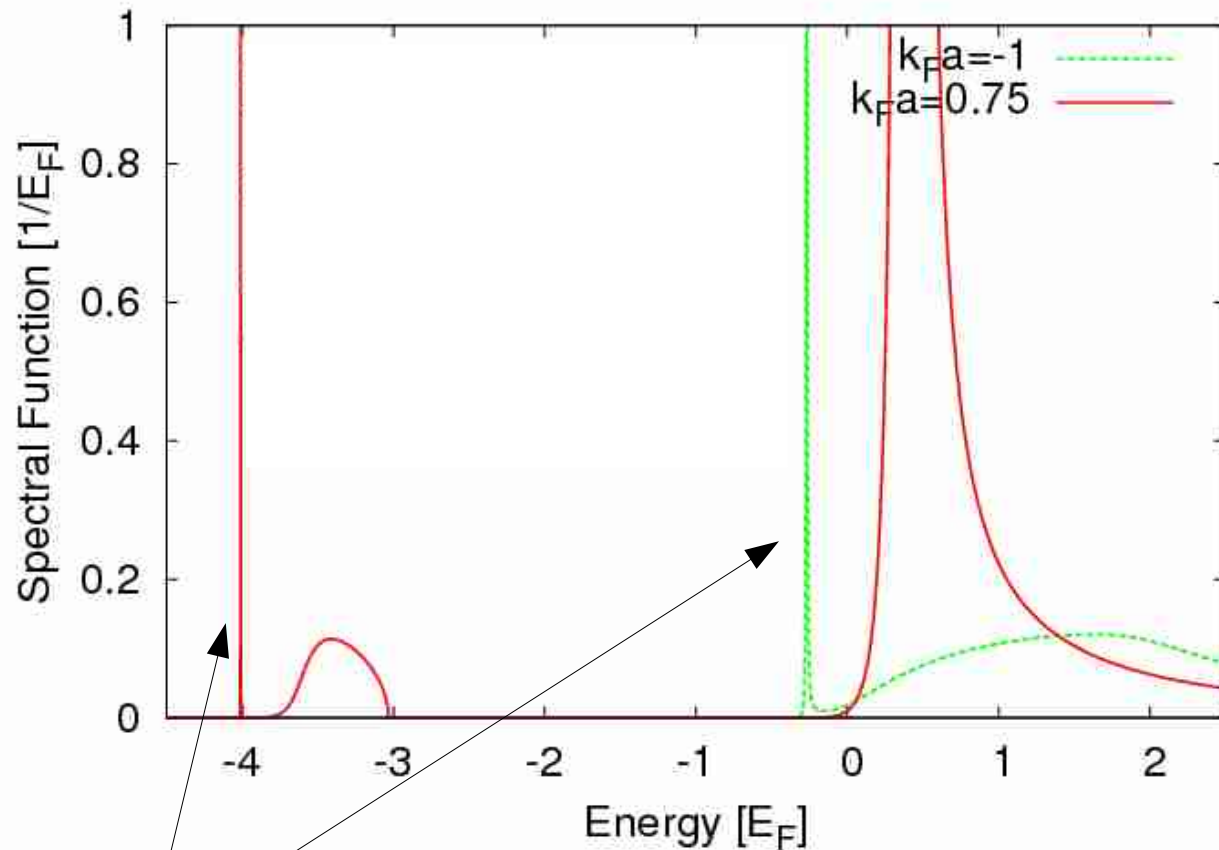
the polaron



the molecule
(heteronuclear?)

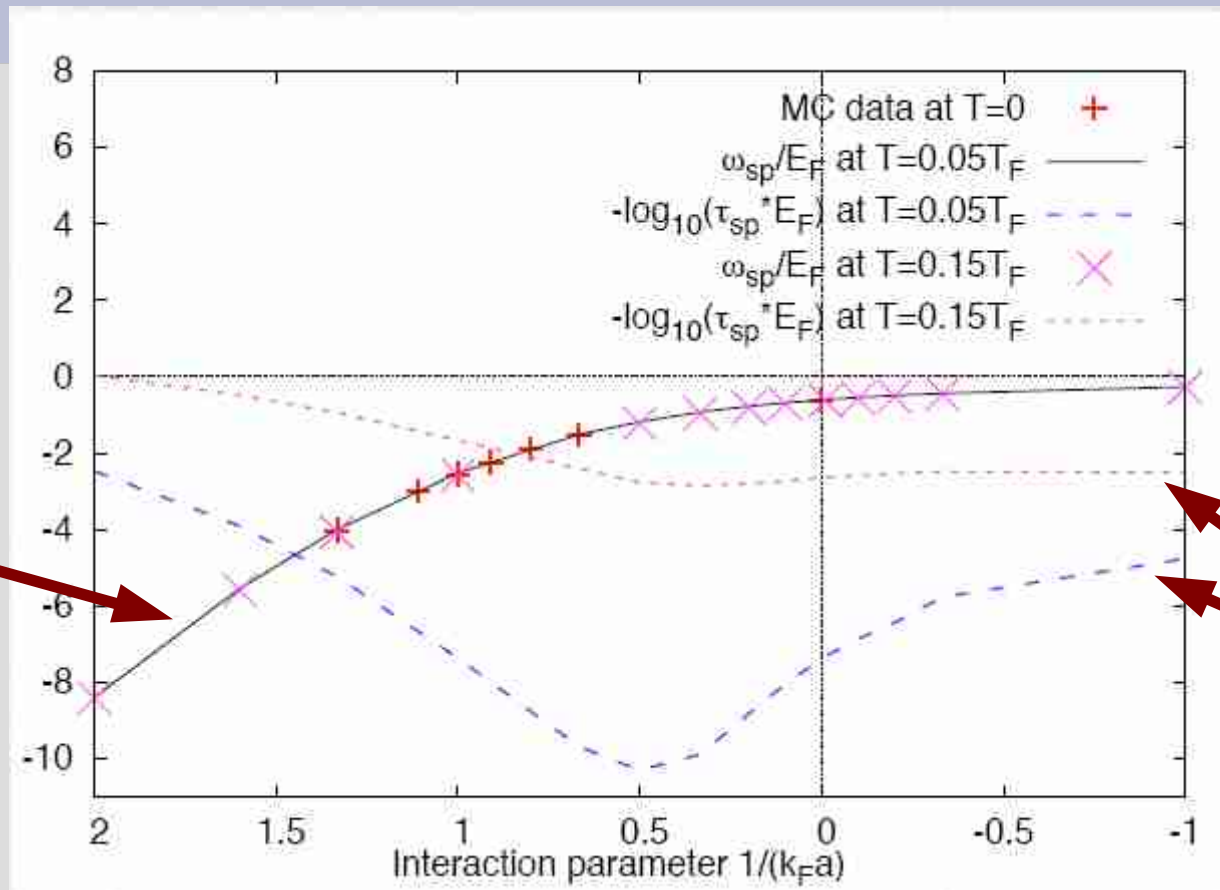


An impurity in a Fermi sea: the spectral function



Narrow feature: a “polaron” (quasiparticle, MB).
It disappears in the BEC limit ($2B$).

Polaron energy and lifetime



Energy

Lifetime

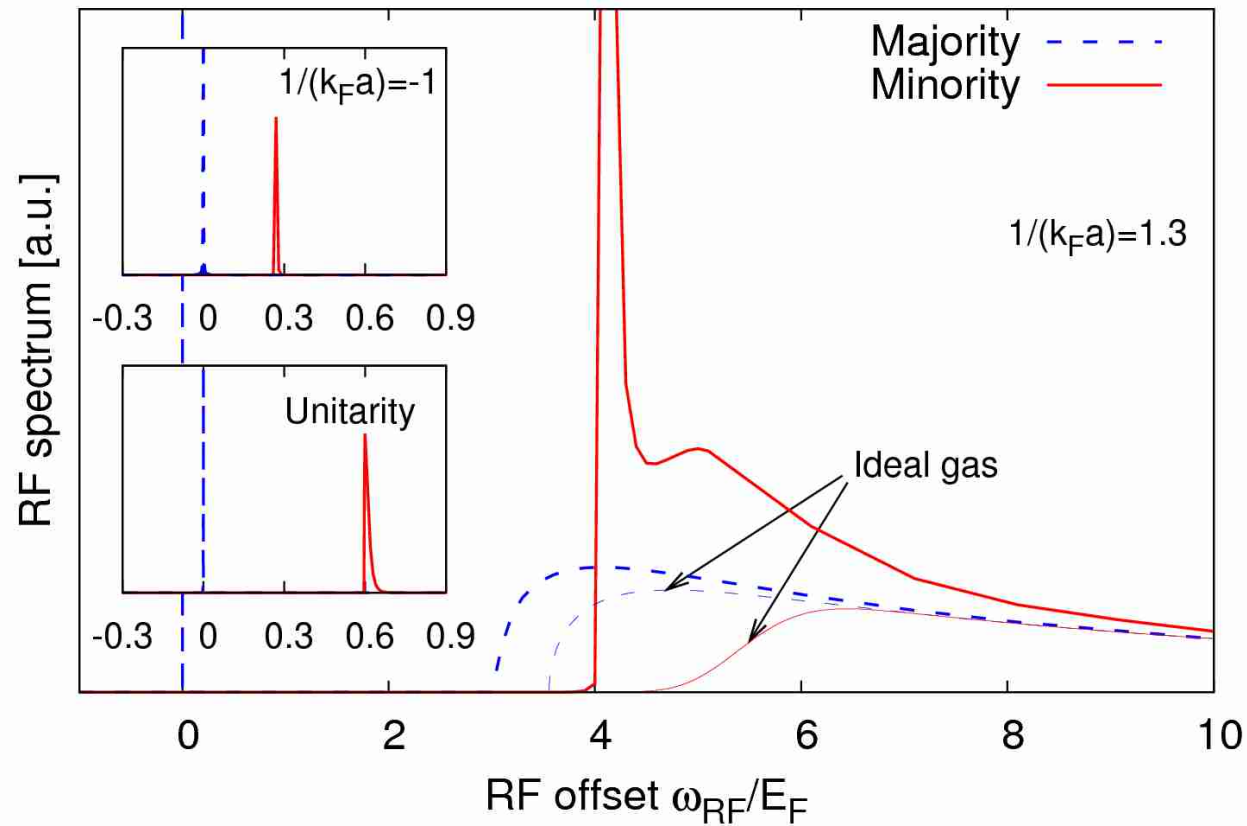
(strong T dependence)

Energy fits with both Monte Carlo and experimental results

(Lobo, Recati, Giorgini and Stringari, PRL 2006; Prokof'ev&Svistunov, PRB 2008; Schirotzek et al., arXiv:0902.3021)

Damping strongly temperature dependent

An impurity in a Fermi sea: the spectrum

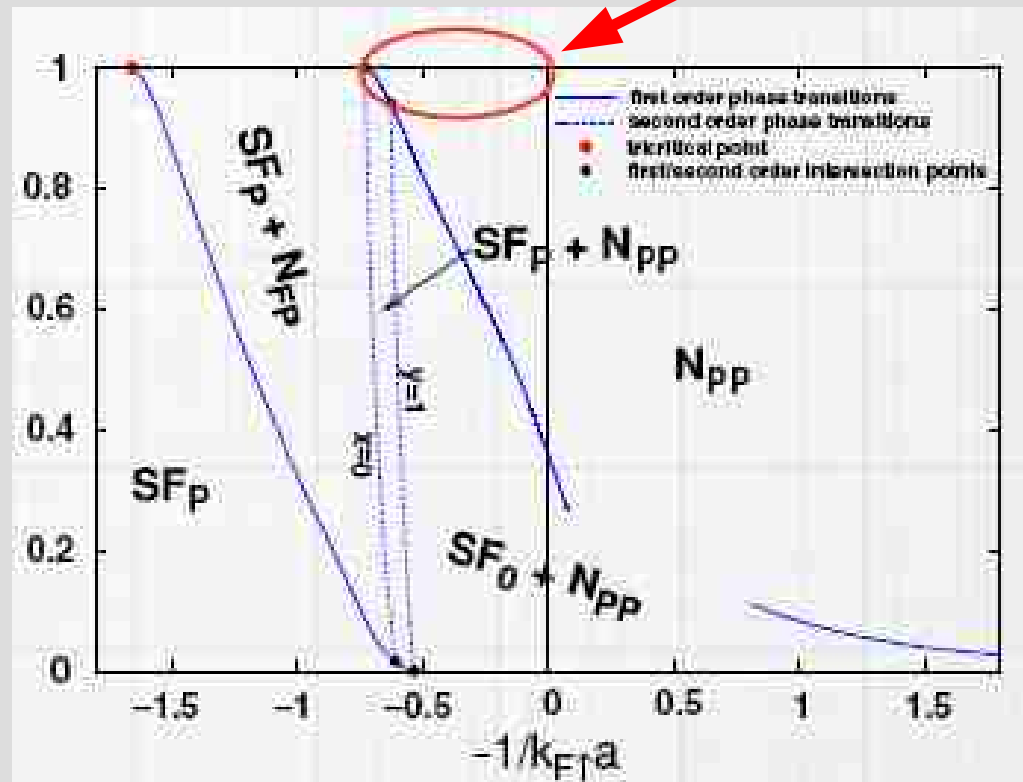


Ideal gas model:
Fermi Golden Rule +
Yukawa molecular function

The spectrum gives info. on both quasiparticle energy and molecular wavefunction

Phase diagram of an imbalanced Fermi gas

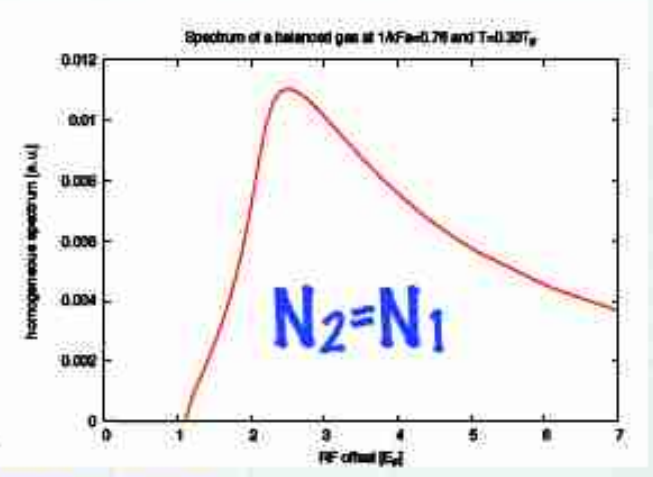
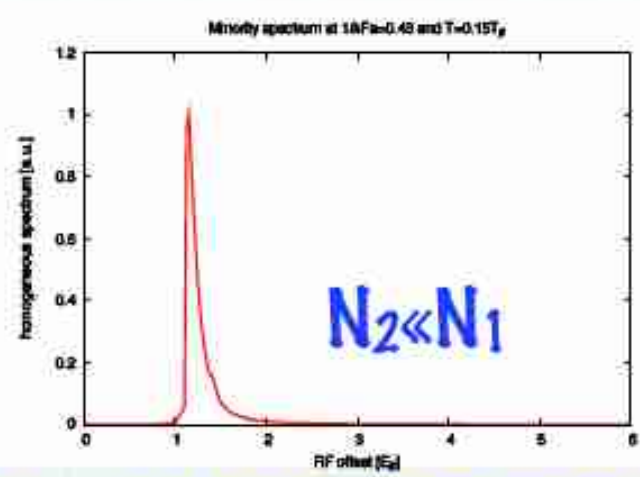
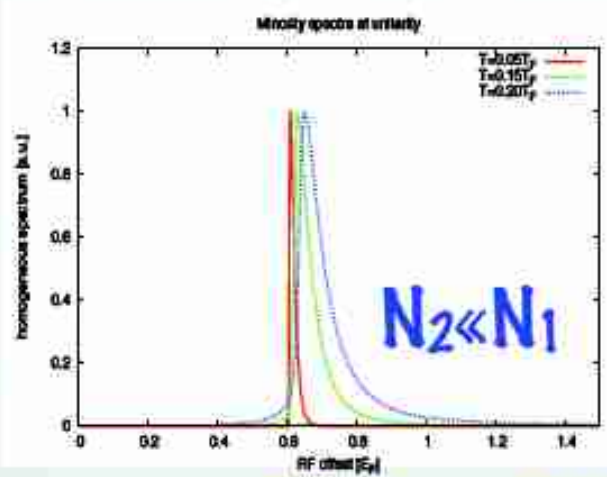
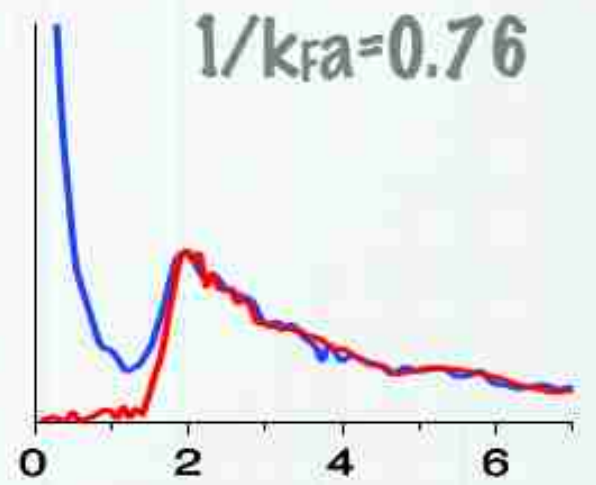
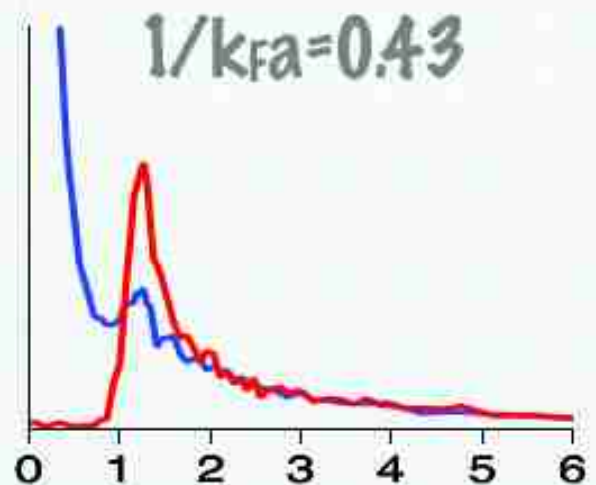
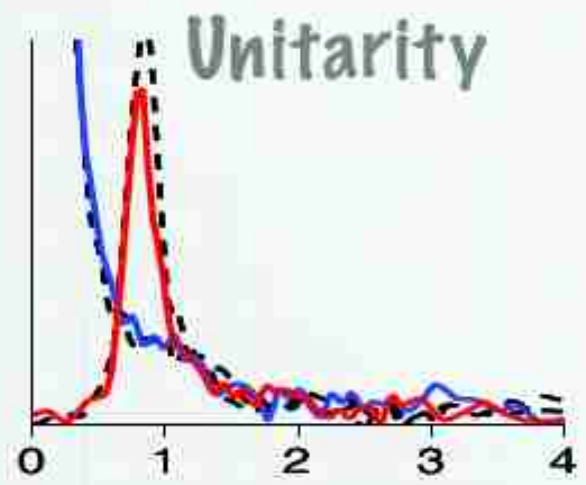
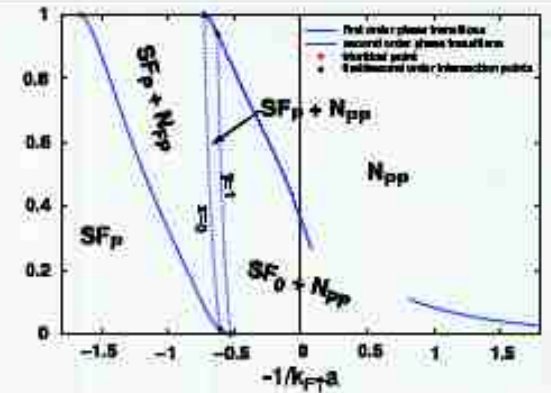
extreme imbalance, strong interactions



S. Pilati & S. Giorgini
PRL 100, 030401 (2008)

RF spectrum

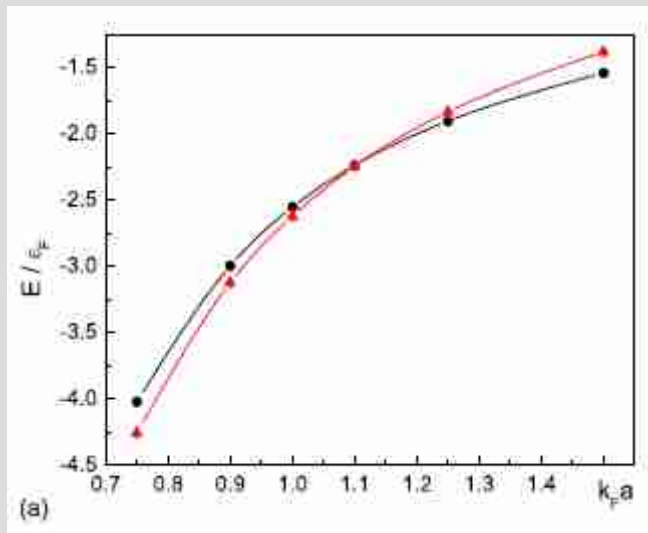
A. Shirotzek et al. arXiv:0902.302



P. Massignan, G. M. Bruun & H. T. C. Stoof PRA 78 031602 (2008)

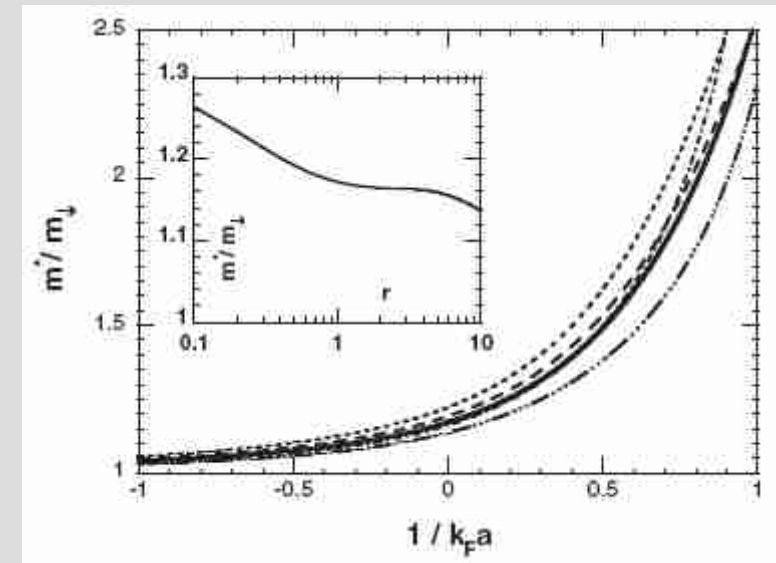
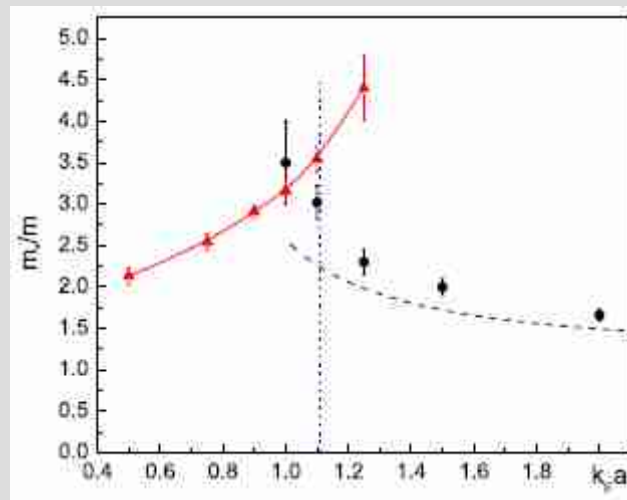
What happens to the impurity on the BEC side?

Energy



(Prokofev & Svistunov, PRB 2008)

Effective mass m^*



(Combescot, Recati, Lobo & Chevy, PRL 2007)

Chevy Ansatz:

$$|\Psi\rangle = \varphi_0 |\mathbf{0}\rangle_{\downarrow} |FS\rangle_{\uparrow} + \sum_{|\mathbf{q}| < k_F < |\mathbf{k}|} \varphi_{\mathbf{k}\mathbf{q}} c_{\mathbf{k}\uparrow}^{\dagger} c_{\mathbf{q}\uparrow} |\mathbf{q} - \mathbf{k}\rangle_{\downarrow} |FS\rangle_{\uparrow}$$

Is it possible to write a single Ansatz able to describe the transition from a polaron to a molecule? (work in progress..)

A **cooking recipe** for: Topological vortex modes and Dirac fermions in a BF mixture

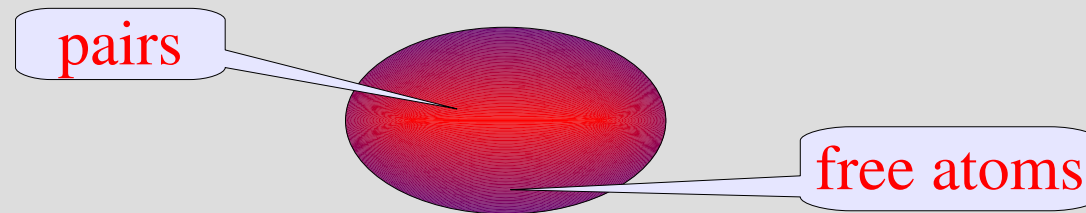
(ongoing work with M. Lewenstein and A. Sanpera)

- **Ingredients:**
 - An optical lattice in 2D
 - 1 boson/site in a Mott state, and some fermions ($0 < F < 1$)
- **Method:**
 - Integrate out the bosons to obtain composite fermions interacting with NN interactions (Lewenstein et al., PRL 2004)
 - A p-wave SF forms at $T < T_c = 10 \sim 100 \text{ nK}$
(stable against 3-body losses via a Quantum Zeno effect, see Han et al., arXiv:0905.2600)
- **Serving suggestions:**
 - This SF supports vortex modes with $E=0$, which have interesting topological properties and non-abelian statistics.
 - At half-filling for the fermions, Dirac fermions appear. **Enjoy!**

(to be published soon, please ask if interested)

Conclusions

- At unitarity there are no real molecules
(double-peaks measured in the spectra come from the trap)



- BEC side ($0 < k_F a < 2$): in a homogeneous gas free atoms coexist locally with stable molecules
- An impurity in a Fermi sea: a polaron (MB) at unitarity, a molecule in the BEC regime (2B).